

# Multiferroics

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[www.lpem.espci.fr/ocg](http://www.lpem.espci.fr/ocg)

# Thanks

*PhD 2011*

R. Schleck

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# Outline

## Introduction to Multiferroics

Definitions

Taxonomy

$\text{BiFeO}_3$  &  $\text{TbMnO}_3$

## Concepts

Symmetry

Ferroelectric transitions

Theoretical digressions

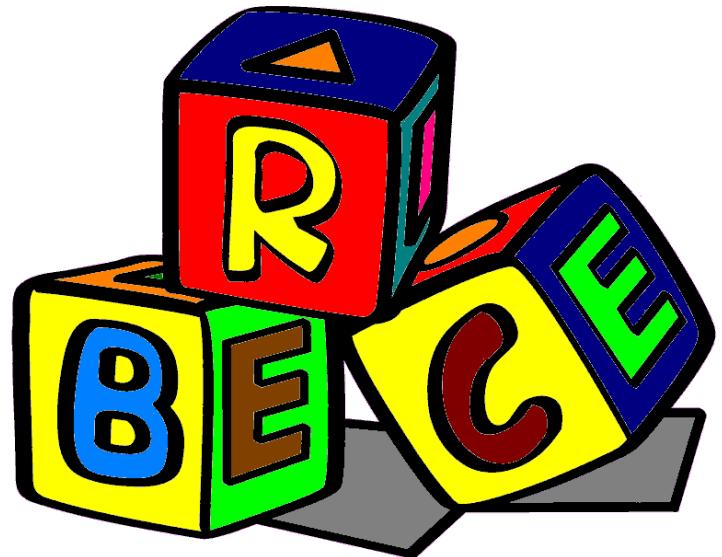
## Optics

Toolbox

Phonons

Electromagnons

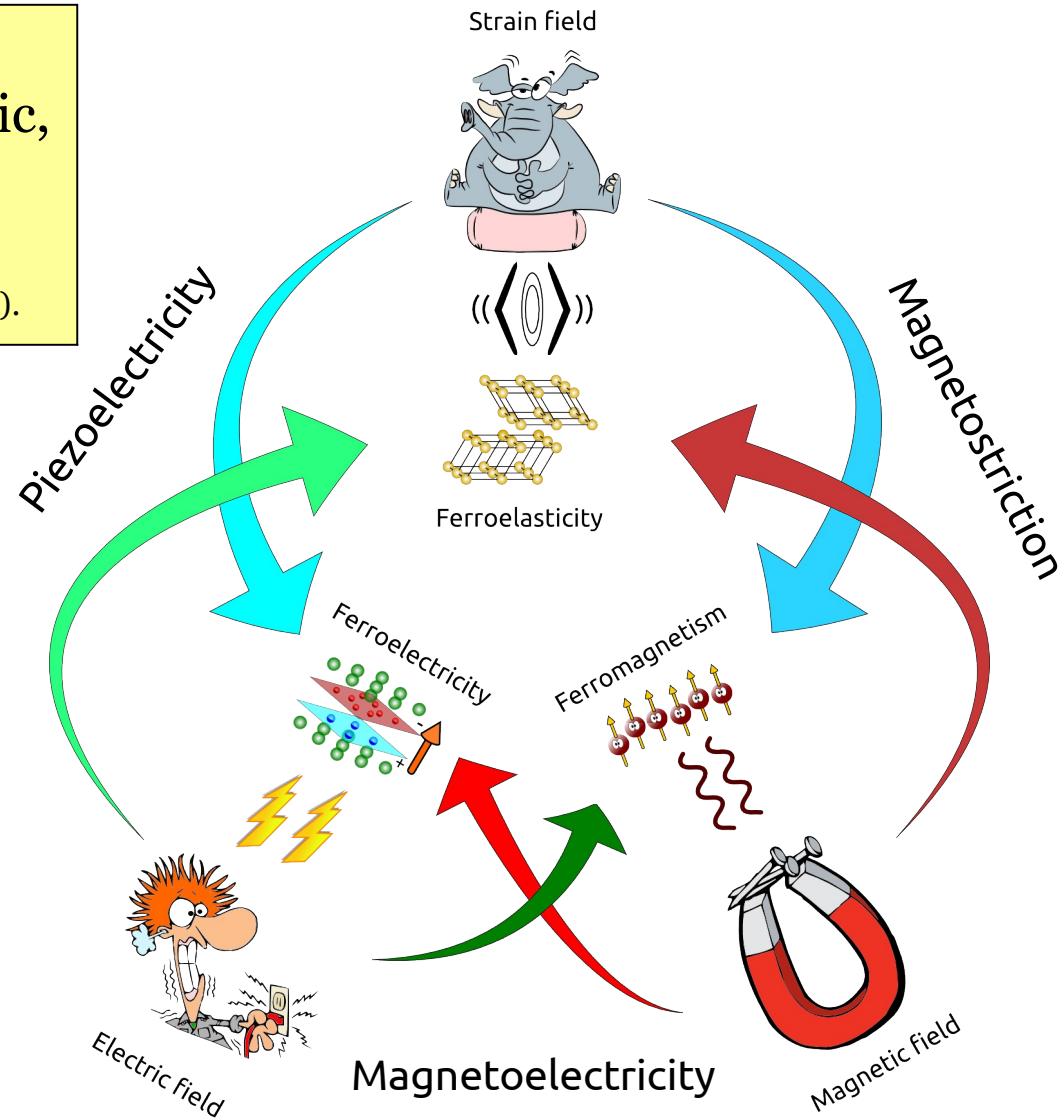
# The Basic Concepts



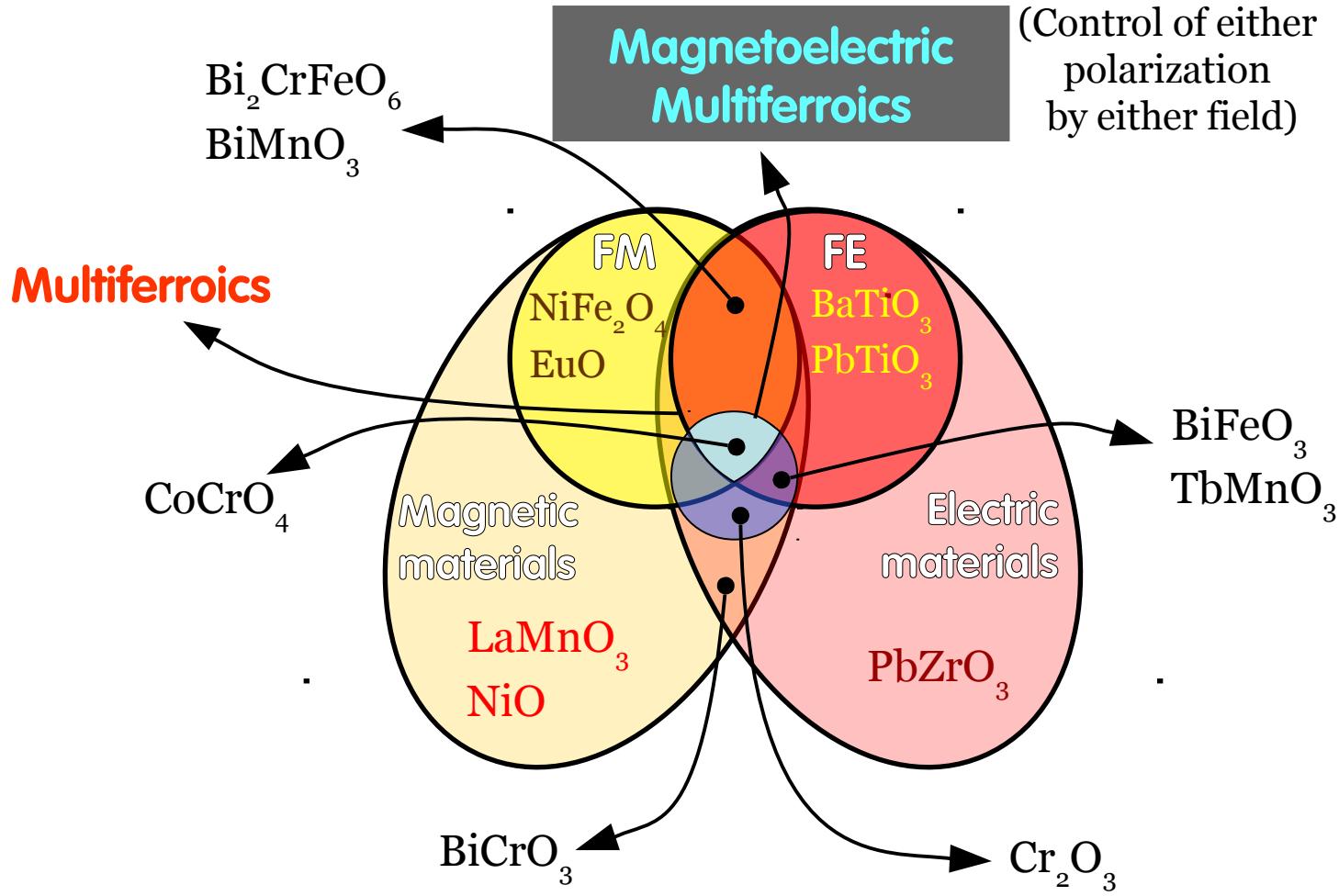
# The Multiferroic Totem

Coexistence of at least two ferroic orders [ferromagnetic, ferroelectric, ferroelastic, ferrotoroidal (??)]

Schmid, Ferroelectrics **162**, 317 (1994).



# The Multiferroic Butterfly



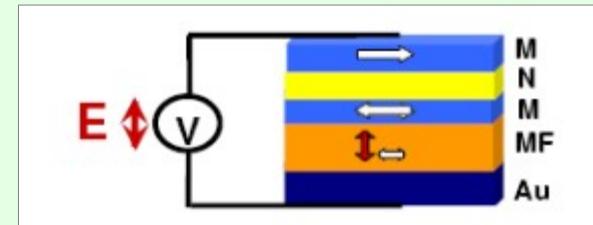
Eerenstein et al., Nature 442, 759 (2006).  
Béa et al, J. Phys. Cond. Matter 20, 434221 (2008)

# Fundraising Excuses

Small coupling – 4 state memory



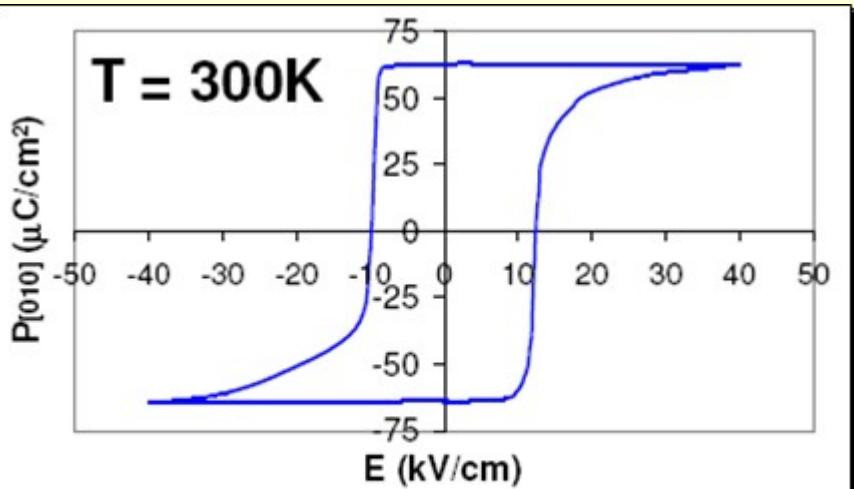
Large coupling – “E” write / “M” read



# Taxonomy of Multiferroics

Ferroelectricity &  
Magnetism coexist

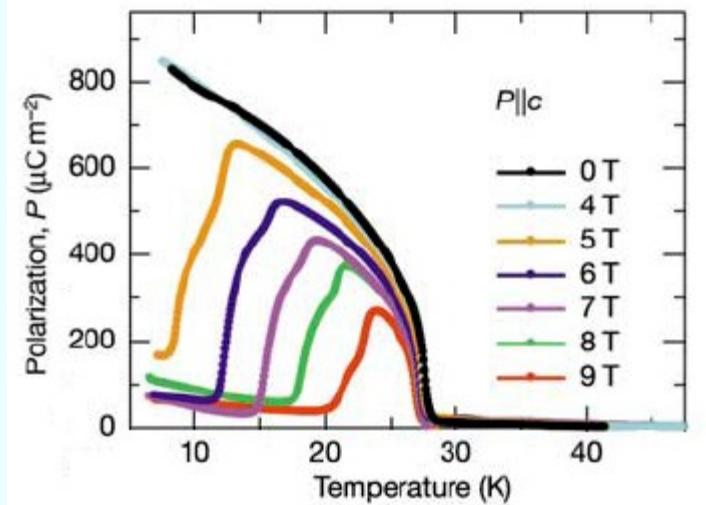
- ✓ BiFeO<sub>3</sub>
- ✓ Large moments
- ✓ High temperature transitions
- ✓ Weak coupling



Lebeugle et al. APL 2007

Magnetism causes  
Ferroelectricity

- ✓ TbMnO<sub>3</sub>
- ✓ Small moments
- ✓ Low temperatures
- ✓ Strong coupling



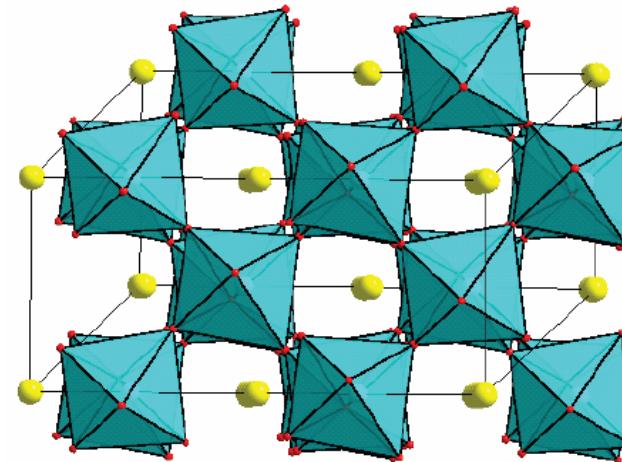
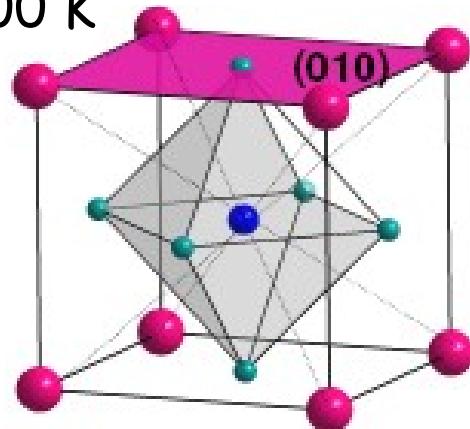
Kimura, Nature 2003

# Orders of Magnitude

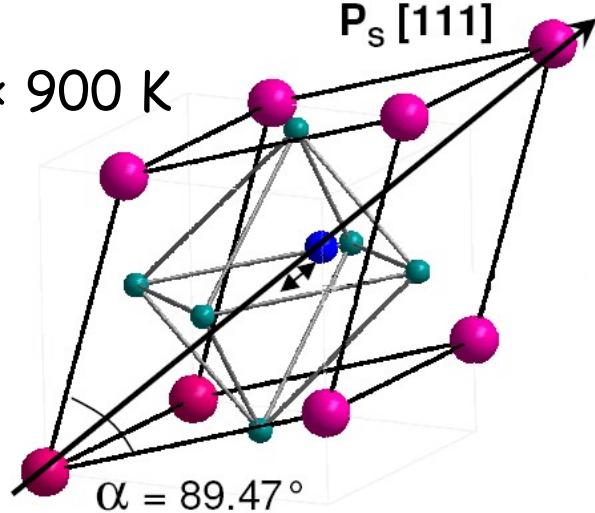
Composition	$T_c$ (K)	P ( $\mu\text{C} / \text{cm}^2$ )	Obs.
$\text{BaTiO}_3$	400	15.	Displacive
$\text{Pb}(\text{Zr},\text{Ti})\text{O}_3$ (PZT)	~600	40.	Displacive
PDVF	(?) ~500	10.	No mag., melts at 450 K
$\text{KH}_2\text{PO}_4$ (KDP)	123	4.	Order-disorder
$\text{K}_2\text{SeO}_4$	93	0.06	Improper FE
$\text{BiFeO}_3$	~1100	100.	MF type I – Lone pairs
$\text{TbMnO}_3$	27	0.08	MF type II
$\text{MnWO}_4$	12	0.005	MF type II

# $\text{BiFeO}_3$ – Ferroelectric / Antiferromagnet

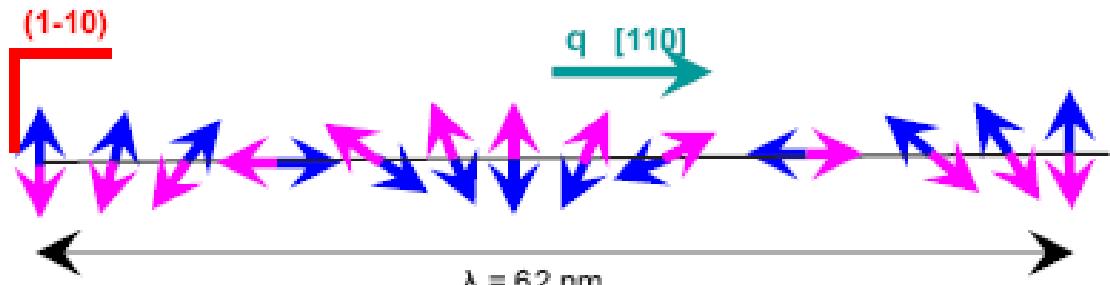
$T > 1100 \text{ K}$



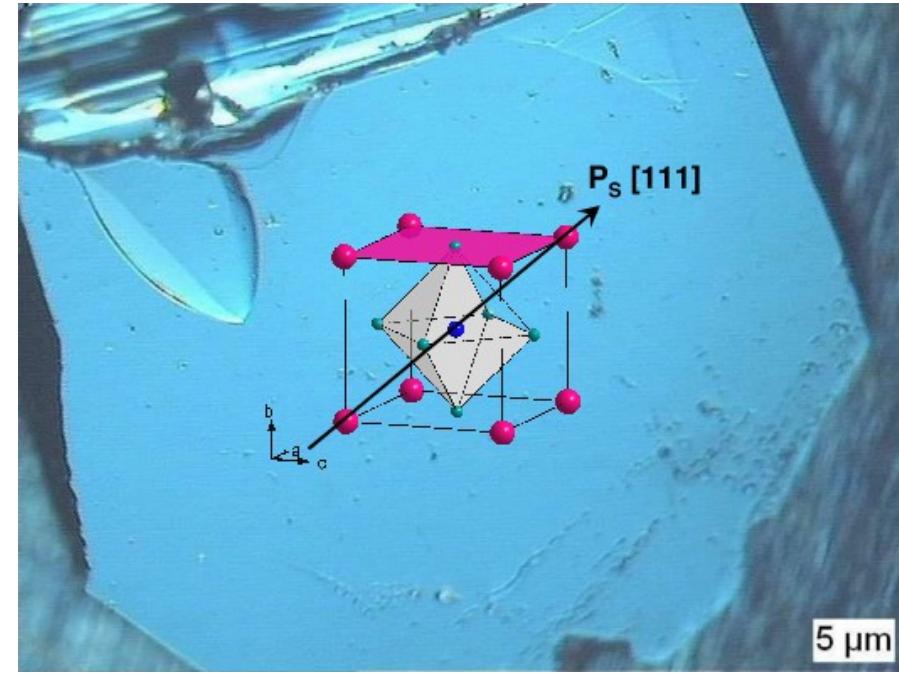
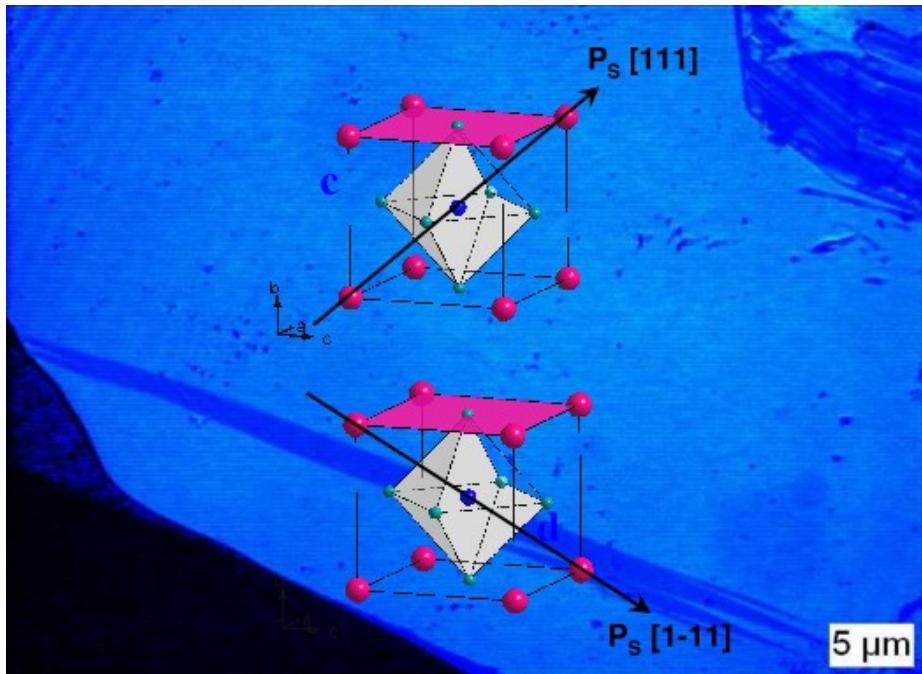
$T < 900 \text{ K}$



- ✓ Rhomboedral distortion of cubic perovskite
- ✓ Incommensurate AF spiral ordering
- ✓  $T_c = 1100 \text{ K} \& T_N = 640 \text{ K}$



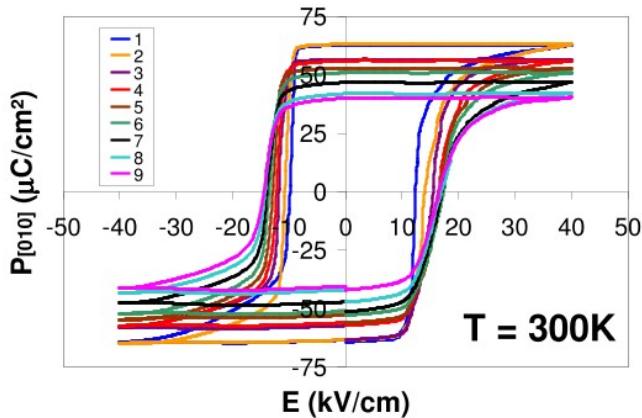
# BiFeO<sub>3</sub> Ferroelastic Domains



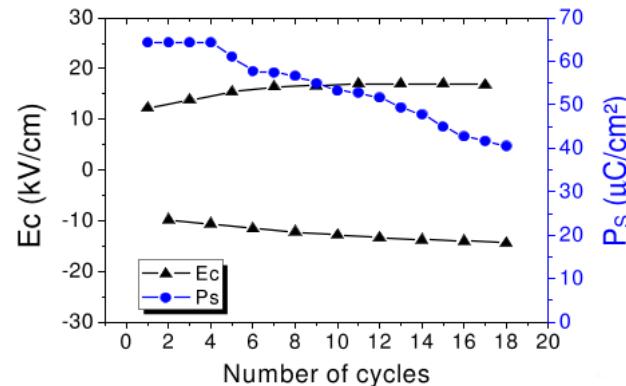
D. Lebeugle, PhD thesis

# $\text{BiFeO}_3$ Properties

Large cyclable polarization

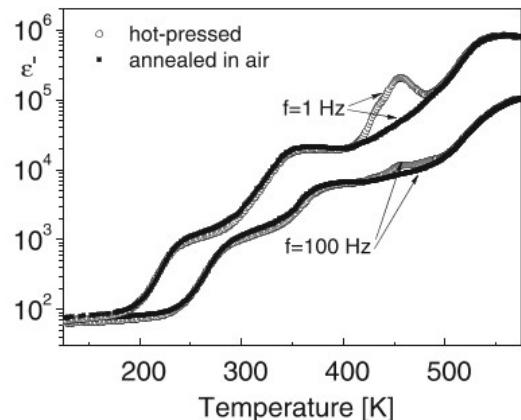


Electric fatigue



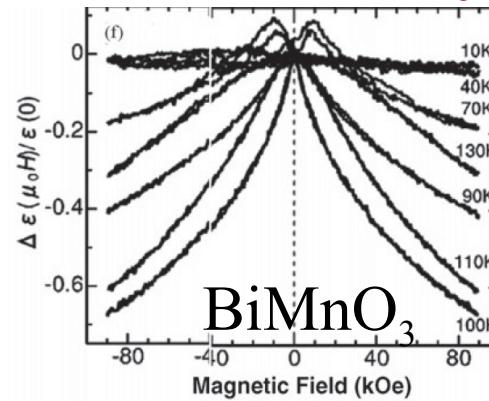
Lebeugle et al., APL **91**, 022907 (2007)

Large dielectric constant



Markiewiks, J. Electroceram **27**, 154 (2011)

Weak coupling ( $\text{BiMnO}_3$ )



T. Kimura PRB **67**, 180401R (2003)

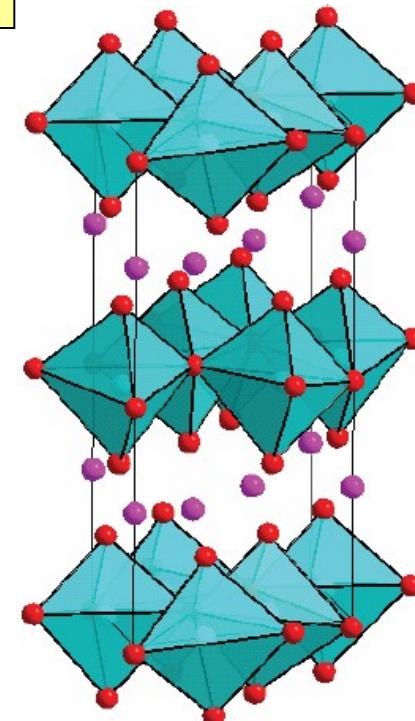
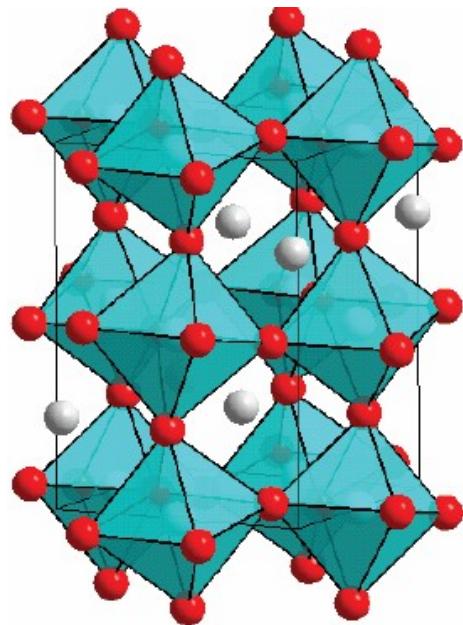
# $\text{RMnO}_3$ structure

Perovskite

La	Ce	Pr	Nd	Pm	Sm	Eu	Gd	Tb	Dy	Ho	Er	Tm	Yb	Lu
----	----	----	----	----	----	----	----	----	----	----	----	----	----	----

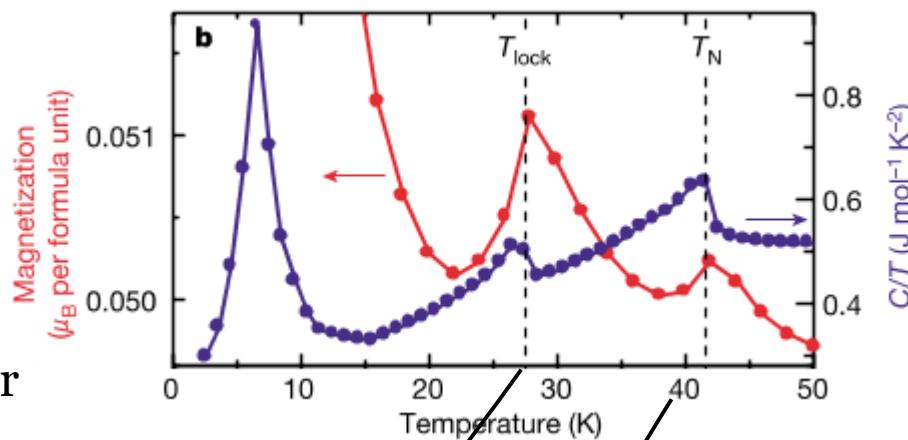
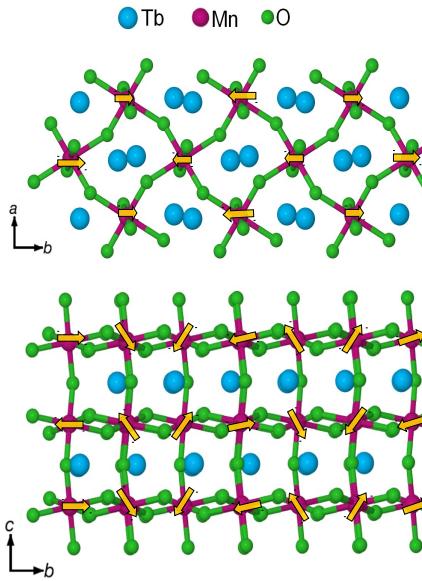
Hexagonal

Y

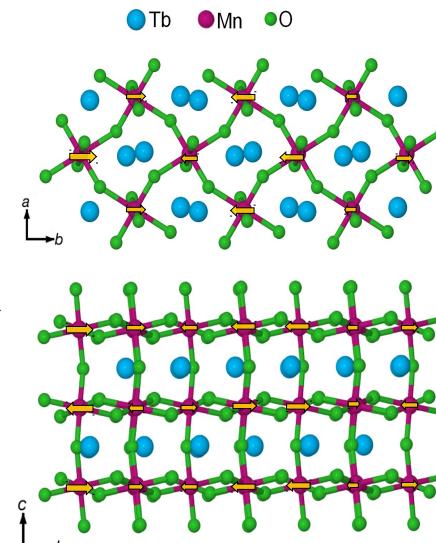


# Multiferroicity in $\text{TbMnO}_3$

$T < 27 \text{ K}$   
Cycloidal spin order  
Ferroelectric

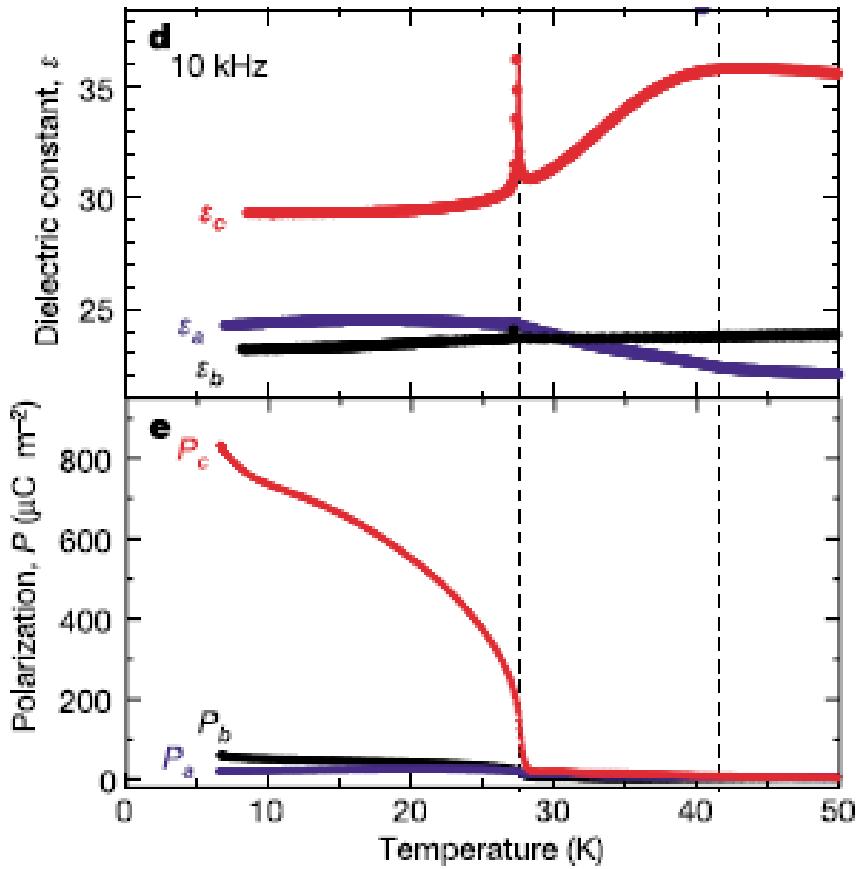


$T < 41 \text{ K}$   
Sinusoidal spin order  
Paraelectric



Kimura *et al*, Nature 426, 55 (2003).

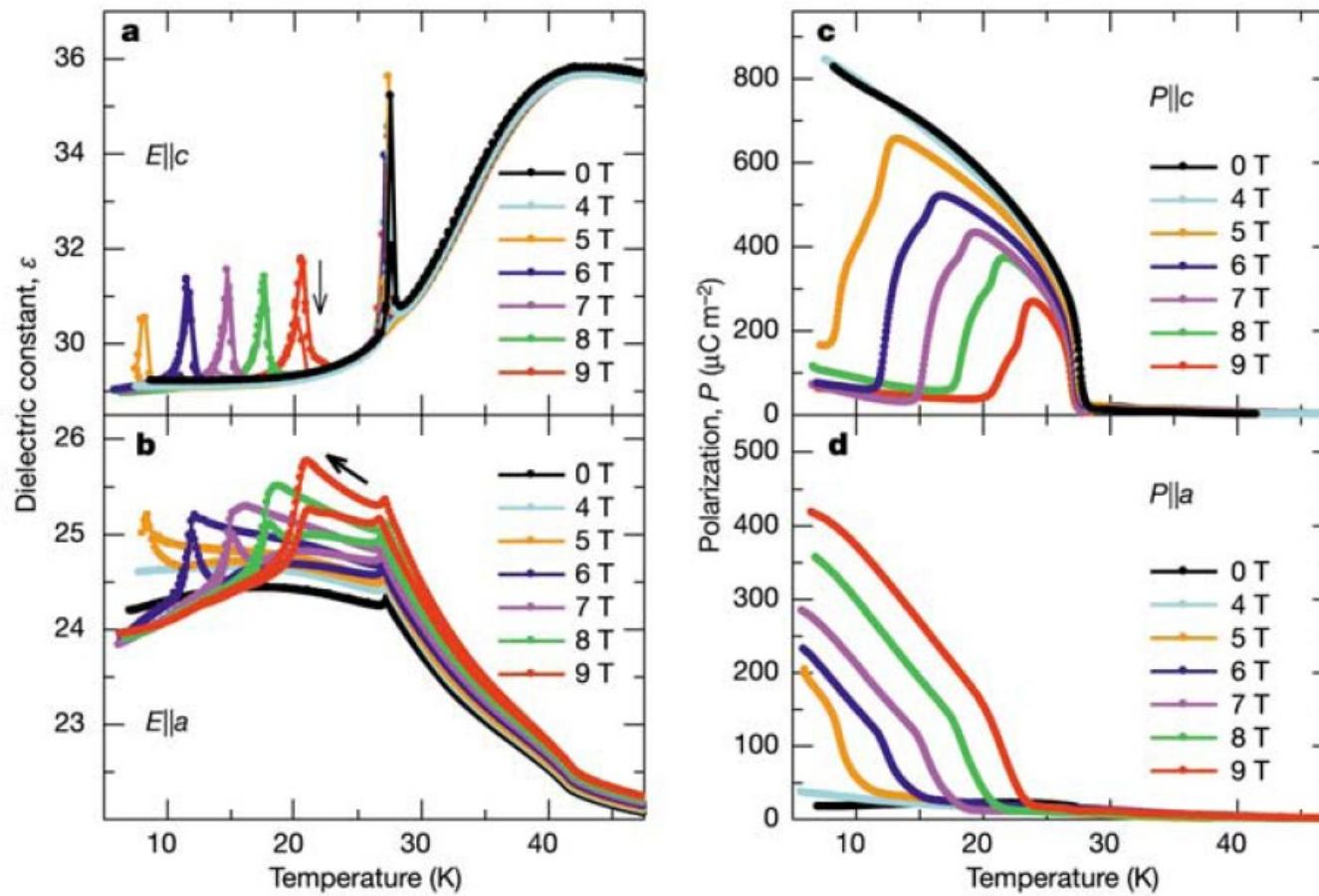
# Ferroelectricity in $\text{TbMnO}_3$



- ✓ Narrow divergence dielectric constant
- ✓ Polarization of the order of that in improper ferroelectrics
- ✓ Ferroelectricity induced by magnetic order

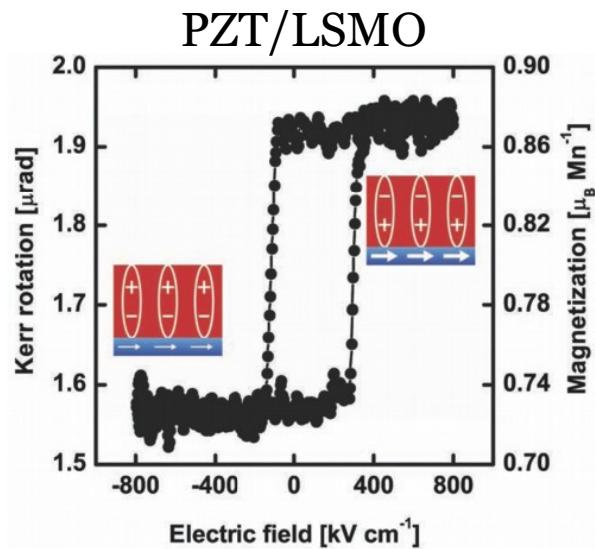
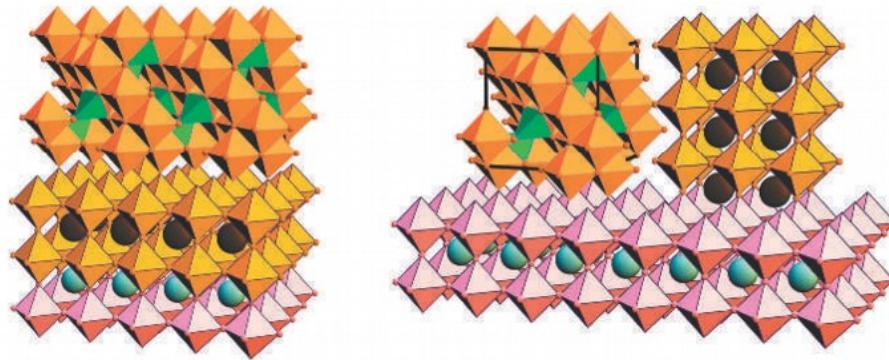
Kimura *et al*, Nature 426, 55 (2003).

# Magneto-electric coupling in $\text{TbMnO}_3$



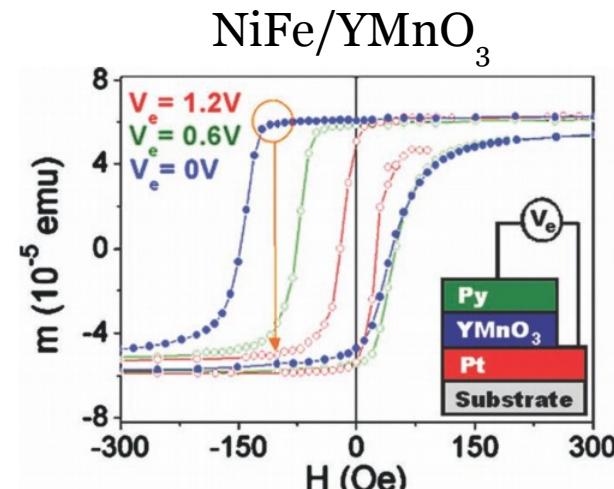
Kimura *et al*, Nature 426, 55 (2003).

# Composite Materials



Molegraaf *et al.*, Ad. Mat. **21**, 3470 (2009)

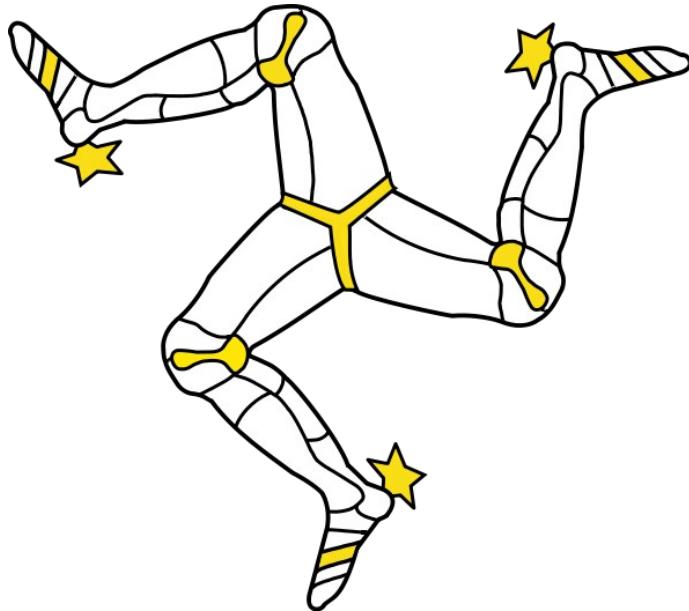
Alternation of piezoelectric and magnetostriictive materials



Laukin, PRL **97**, 227201 (2006)

Fiebig, J. Appl. Phys. D **38**, R123 (2005)  
Ramesh and Spalding, Nat. Mat. **6**, 22 (2007)  
Vaz et al. Ad. Mat. **22**, 2900 (2010)

# Symmetry Considerations

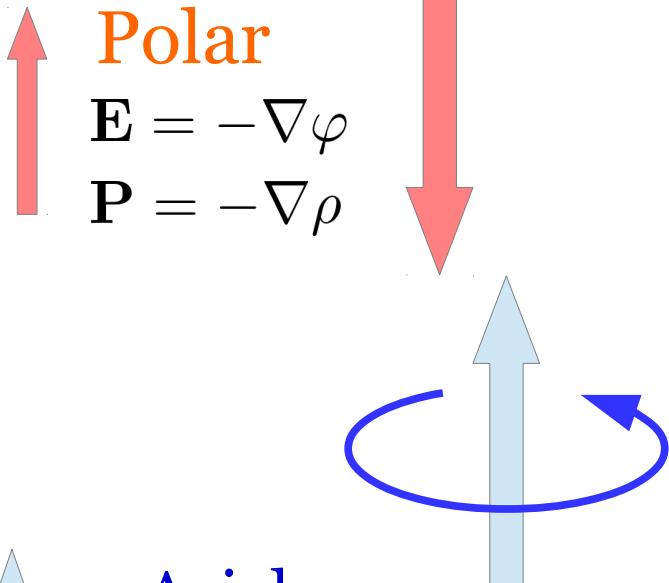


# Vectors and Symmetry Operations

Polar

$$\mathbf{E} = -\nabla\varphi$$

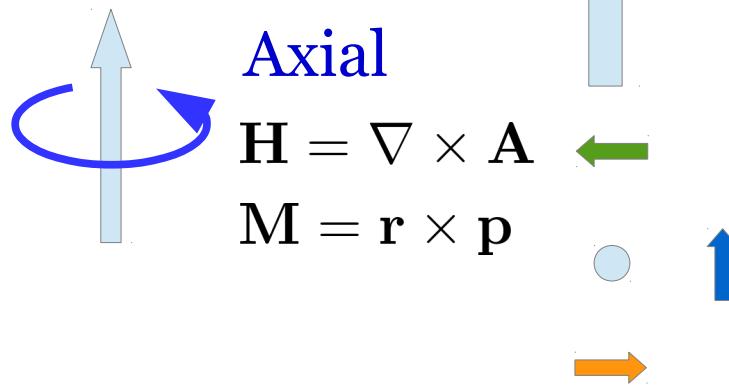
$$\mathbf{P} = -\nabla\rho$$



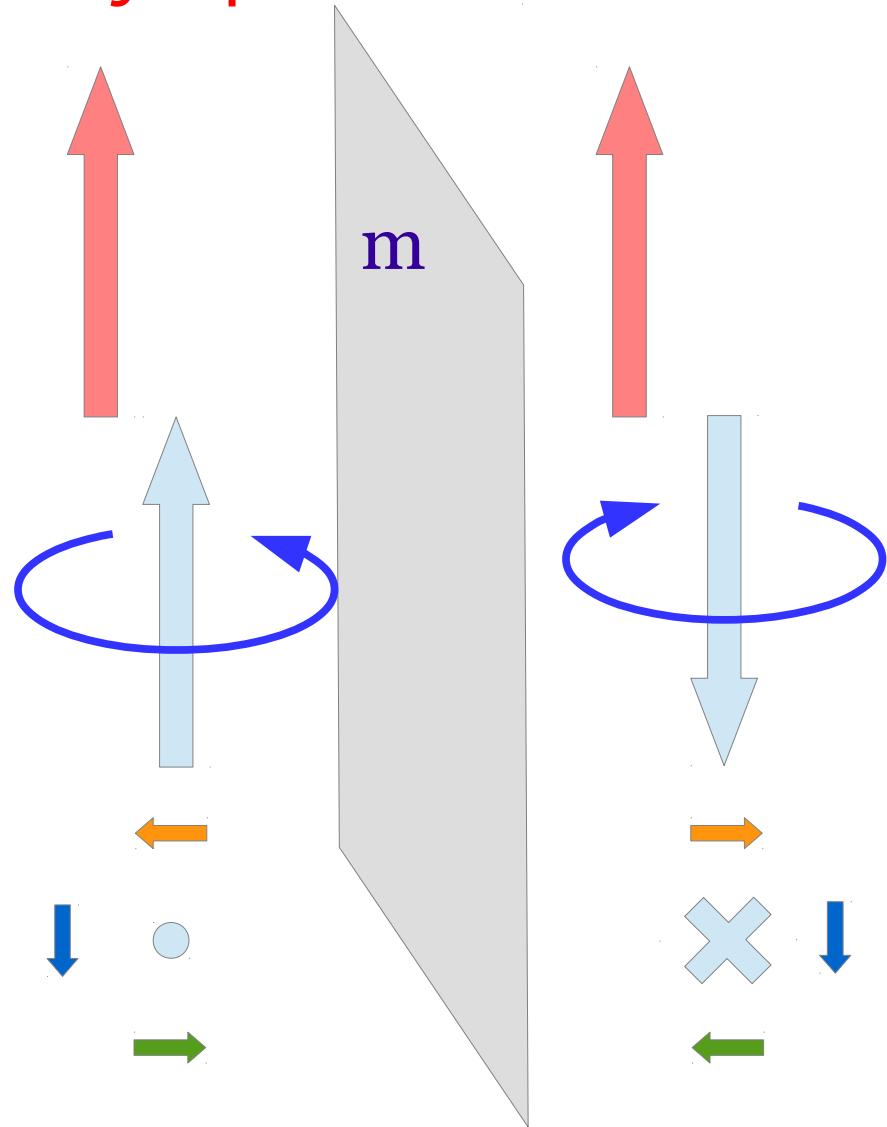
Axial

$$\mathbf{H} = \nabla \times \mathbf{A}$$

$$\mathbf{M} = \mathbf{r} \times \mathbf{p}$$



Inv.



# The first ideas...

CURIE. — SYMÉTRIE DANS LES PHÉNOMÈNES PHYSIQUES. 393

## SUR LA SYMÉTRIE DANS LES PHÉNOMÈNES PHYSIQUES, SYMÉTRIE D'UN CHAMP ÉLECTRIQUE ET D'UN CHAMP MAGNÉTIQUE;

PAR M. P. CURIE.

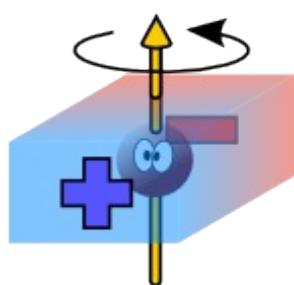
*J. de Phys.*, 3<sup>e</sup> série, t. III. (Septembre 1894.)

Same year as Rochelle salt (Pockels)

Les phénomènes généraux de l'électricité et du magnétisme nous indiquent donc seulement une *liaison entre les symétries* du champ électrique et du champ magnétique, de telle sorte que, si l'on adopte (*c*) pour la symétrie de l'un, il faut admettre (*d*) pour la symétrie de l'autre, et réciproquement. Pour lever cette indétermination, il faut faire intervenir d'autres phénomènes, les phénomènes électrochimiques ou d'électricité de contact, les phénomènes pyro ou piézoélectriques, ou encore le phénomène de Hall, ou celui de la polarisation rotatoire magnétique.

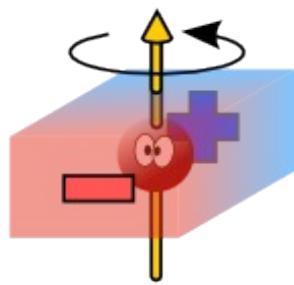
# Symmetry Shattering

+m & +s



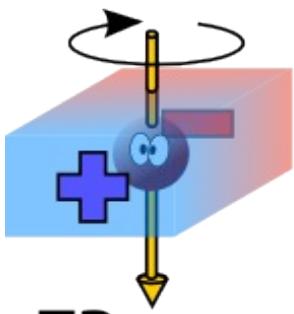
SI:

-m & +s

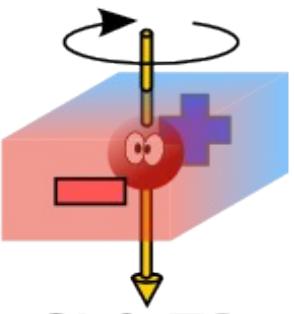


TR:

+m & -s



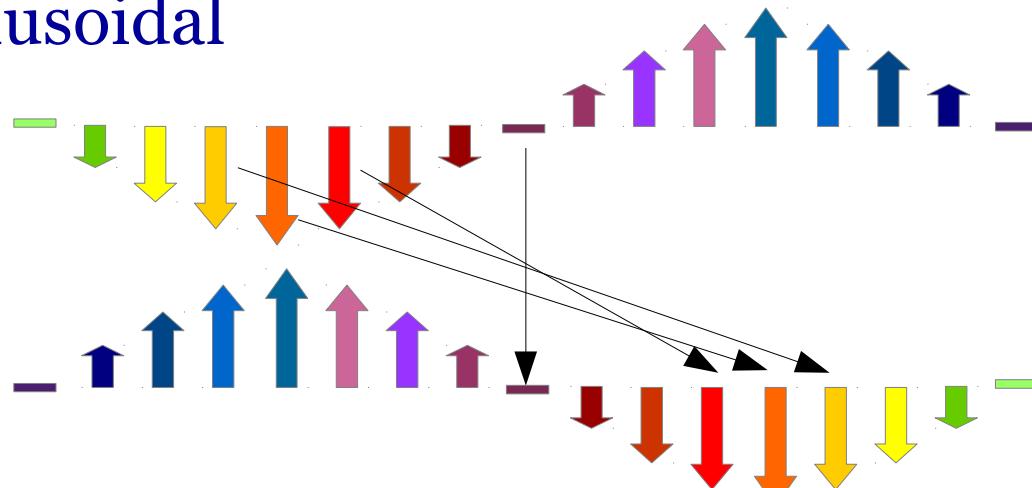
SI & TR:  
-m & -s



- ✓ Spatial Inversion (SI) symmetry breaking:  
*Makes sure that the electrical polarization exists but has no effect on the magnetization*
- ✓ Time Reversal (TR) symmetry breaking:  
*Allows for spontaneous magnetic ordering but ignores polar vector quantities*
- ✓ SI + TR breaking  
*Both, spontaneous **P** and **M** exist*

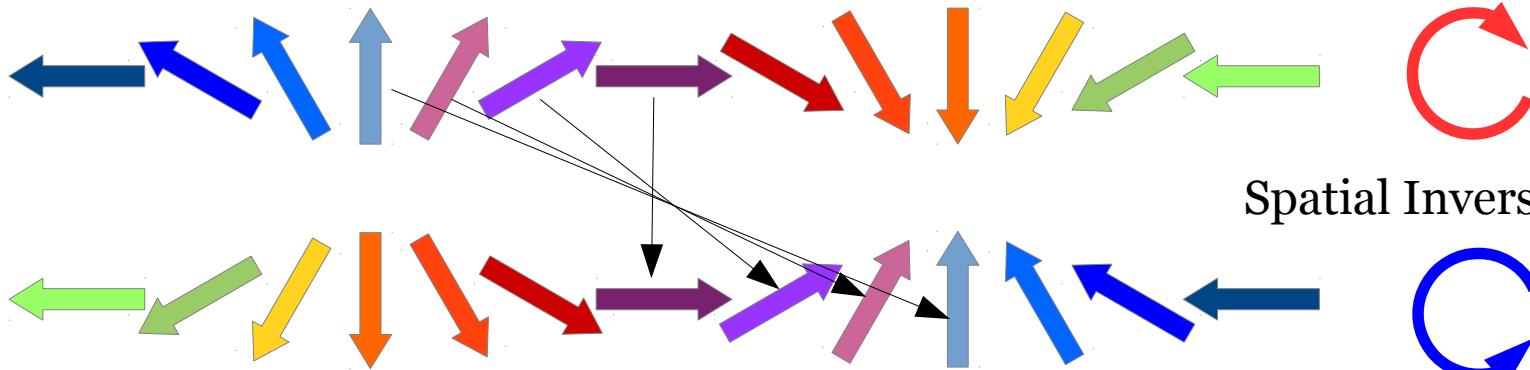
# Spin Ordering

Sinusoidal



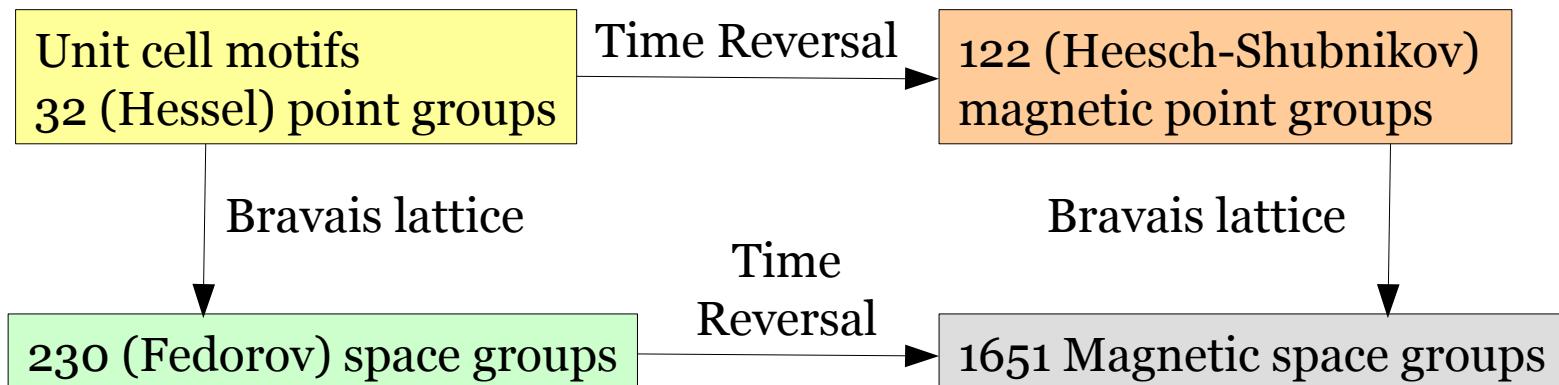
No Spatial Inversion breaking

Cycloidal



Spatial Inversion broken

# Symmetry groups



Magnetic point group	Required relations between $P_s$ and $M_s$	Ferroelectrically induced spin canting
1	none	allowed
$m'$	$\hat{P} \parallel m', \hat{M} \parallel m'$	allowed
$m$	$\hat{P} \parallel m, \hat{M} \perp m$	allowed
$2'$	$\hat{P} \perp \hat{M}, \hat{P} \parallel 2'$	allowed
$m'm'2'$	$\hat{P} \parallel 2', \hat{M} \perp m$	allowed
$2, m'm'2$	$\hat{P} \parallel 2 \parallel \hat{M}$	forbidden
$3, 3m', 4, 4m'm', 6, 6m'm'$	$\hat{P} \parallel \hat{M} \parallel$ principal axis	forbidden

112 H-S groups  
31 allow **P**  
31 allow **M**  
13 allow **P & M**

N. Hill (Spaldin), J. Phys. Chem. B **104**, 6694 (2000)

# Magnetoelectric Effect

$$\begin{aligned} -g(\mathbf{E}, \mathbf{H}) = & \cdots + P_i E_i + M_i H_i \\ & + \frac{1}{2} \varepsilon_0 \varepsilon_{ik} E_i E_k + \frac{1}{2} \mu_0 \mu_{ik} H_i H_k \\ & + \alpha_{ik} E_i H_k \quad \xrightarrow{\hspace{10em}} \text{Linear coupling} \\ & + \frac{1}{2} \beta_{ijk} E_i H_j H_k + \frac{1}{2} \gamma_{ijk} H_i E_j E_k \\ & + \frac{1}{6} \delta_{ijk} E_i E_j E_k + \frac{1}{6} \eta_{ijk} H_i H_j H_k + \cdots \end{aligned} \quad \left. \right\} \text{Bilinear coupling}$$

$$\alpha_{ik} \sim 10^6 \beta_{ijk}$$

$$\begin{aligned} P_k(\mathbf{E}, \mathbf{H}) = & -\frac{\partial g}{\partial E_k} = \cdots + P_k + \varepsilon_0 \varepsilon_{ki} E_i + \alpha_{ki} H_i \\ & + \frac{1}{2} \beta_{kij} H_i H_j + \gamma_{ijk} H_i E_j + \cdots \\ M_k(\mathbf{E}, \mathbf{H}) = & -\frac{\partial g}{\partial H_k} = \cdots + M_k + \mu_0 \mu_{ki} H_i + \alpha_{ik} E_i \\ & + \beta_{ijk} E_i H_j + \frac{1}{2} \gamma_{kij} E_i E_j + \cdots \end{aligned}$$

Rivera, Eur. Phys. J. B 71, 299 (2009)

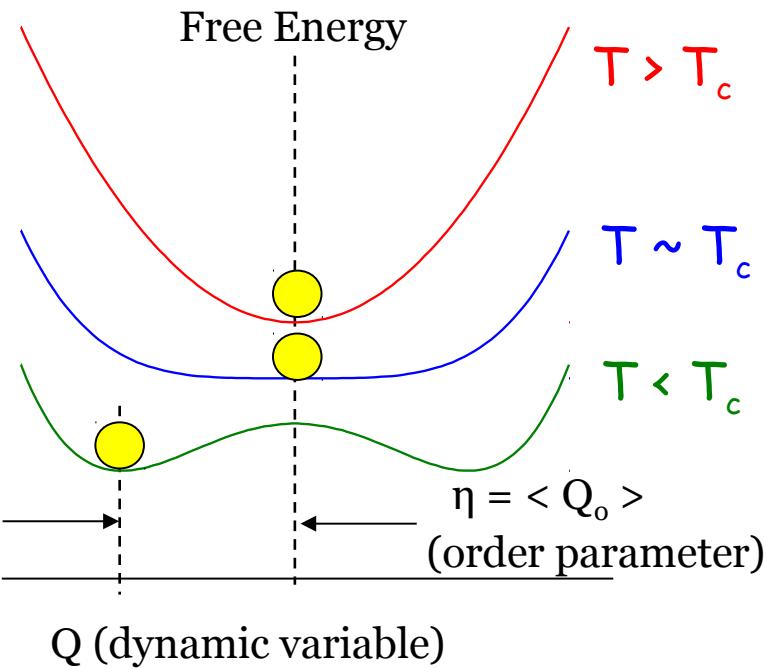
# Strong Ferroelectric Phase Transitions



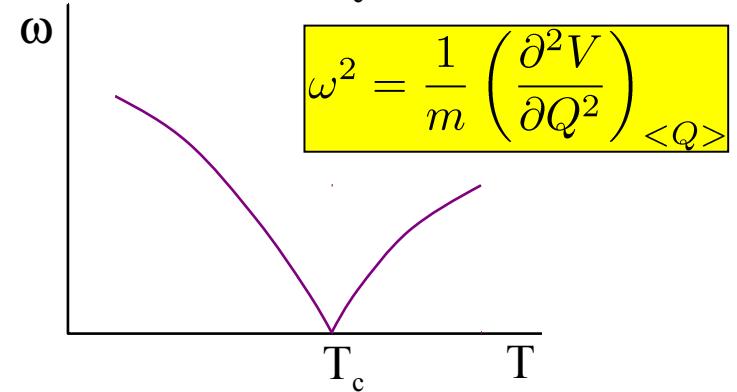
# Soft Mode & Phase Transitions

$Q$  is a dynamic variable of the system

Order parameter  $\eta = \langle Q \rangle$   $\begin{cases} = 0, T > T_c \\ \neq 0, T < T_c \end{cases}$

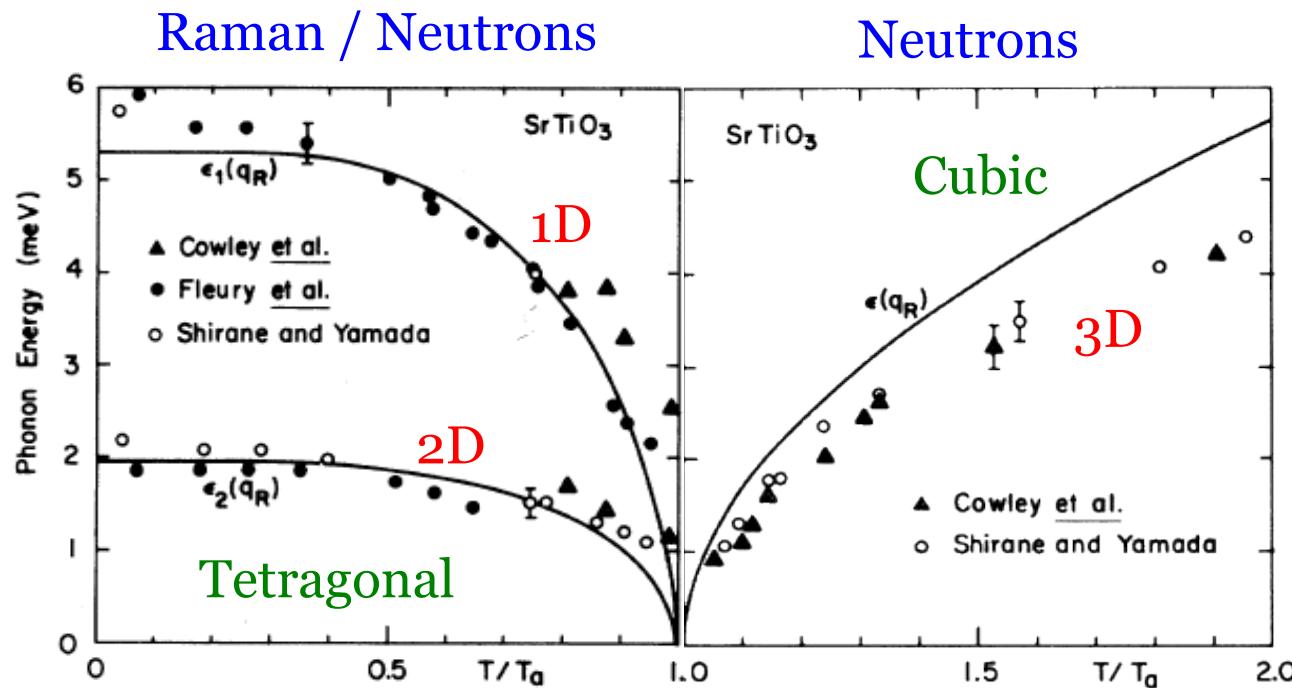


$$\langle Q \rangle \quad \left( \frac{\partial V}{\partial Q} \right)_{\langle Q \rangle} = 0$$



# Is this soft mode mambo-jambo real?

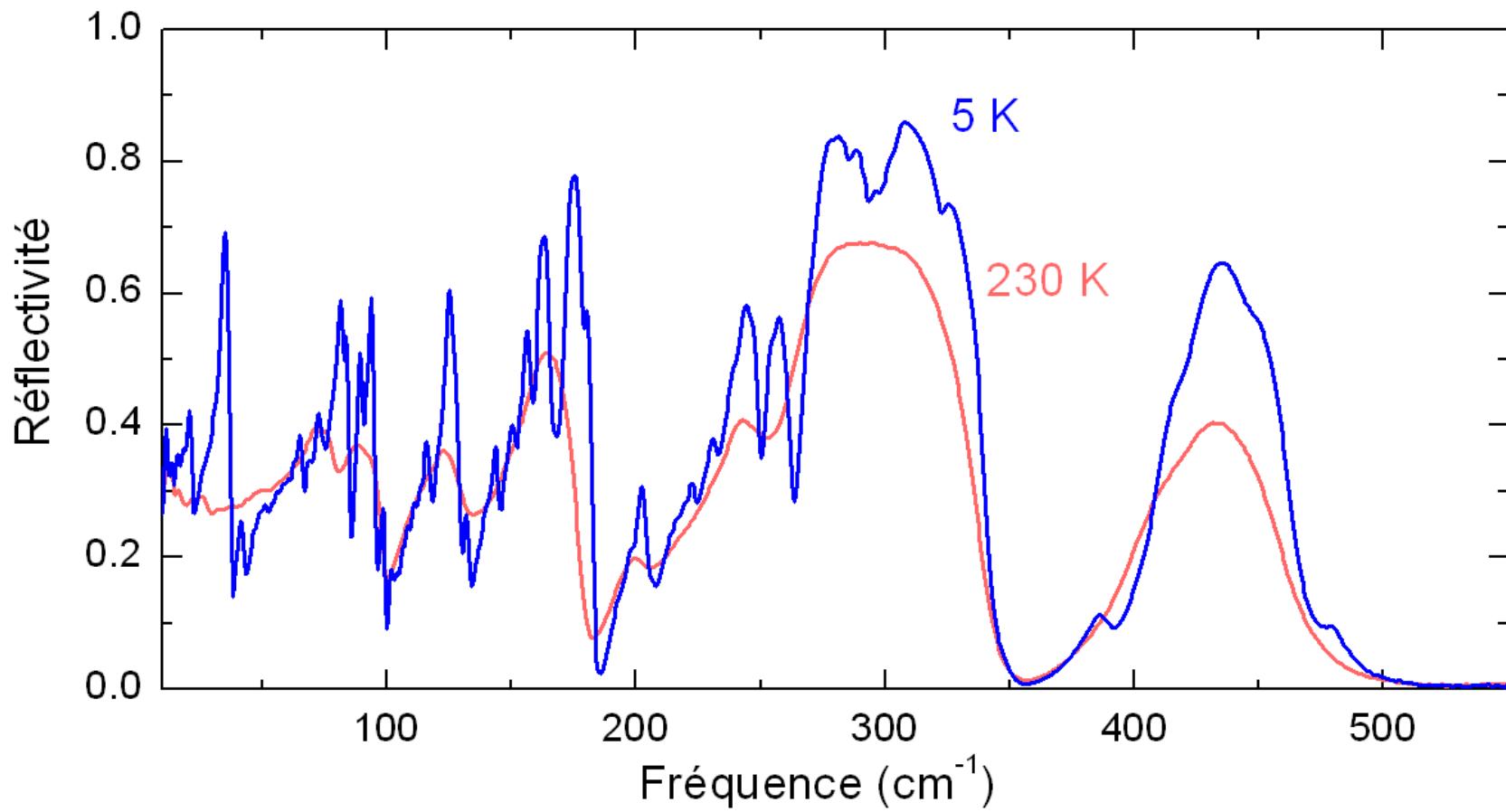
$\text{SrTiO}_3$ : Cubic  $\rightarrow$  Tetragonal transition at  $T_c = 110 \text{ K}$



Soft mode = Phonon  $\rightarrow$  Displacive transition

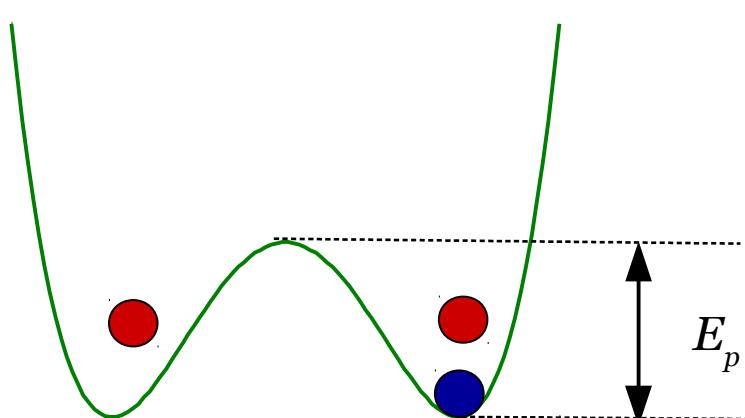
Feder & Pytte, PRB 1, 4803 (1970)

# Orthorhombic-triclinic(?) transition in $\text{BaMnF}_4$



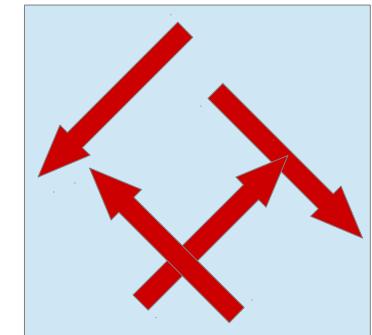
Schleck et al. 2012.

# Order-Disorder transitions



●  $k_B T > E_p$

Random polarization orientation

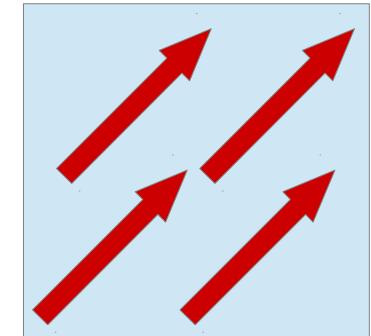


$$k_B T_c \sim E_p$$

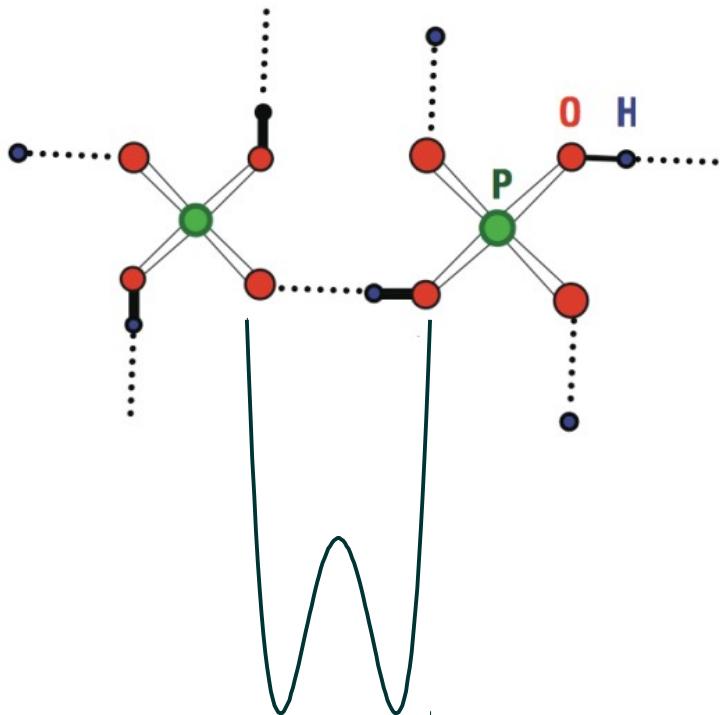
- Same potential for paraelectric and ferroelectric phases
- “Soft-mode” with a relaxation character

●  $k_B T < E_p$

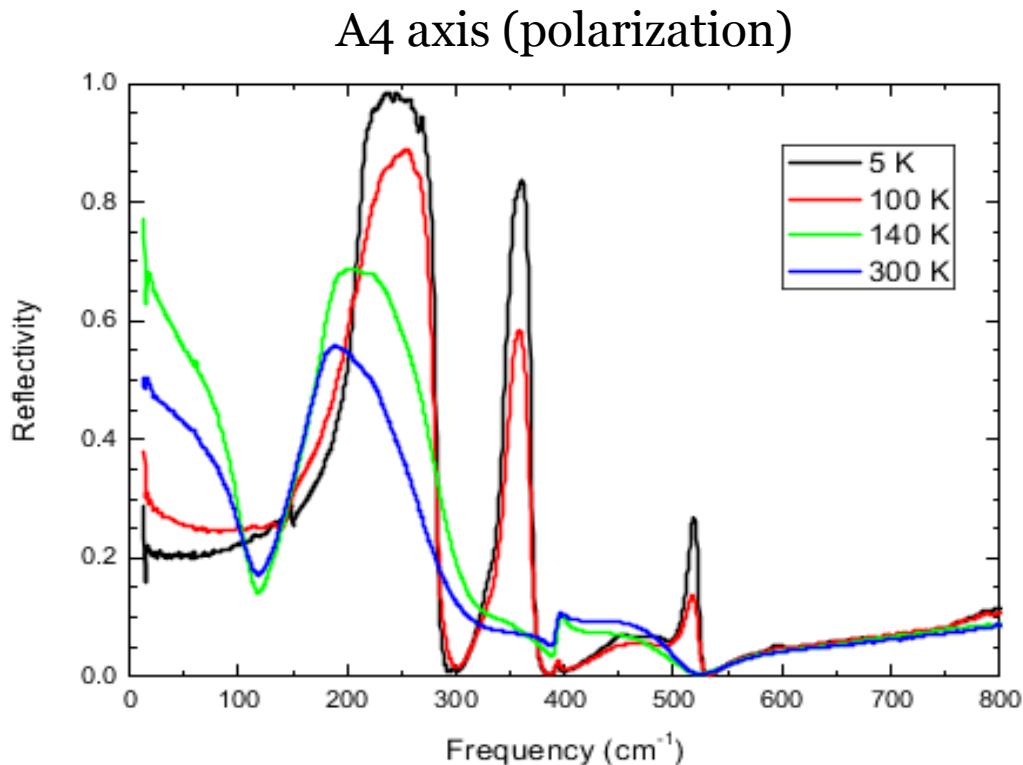
One polarization direction freezes



# Ferroelectricity in $\text{KH}_2\text{PO}_4$ (KDP)

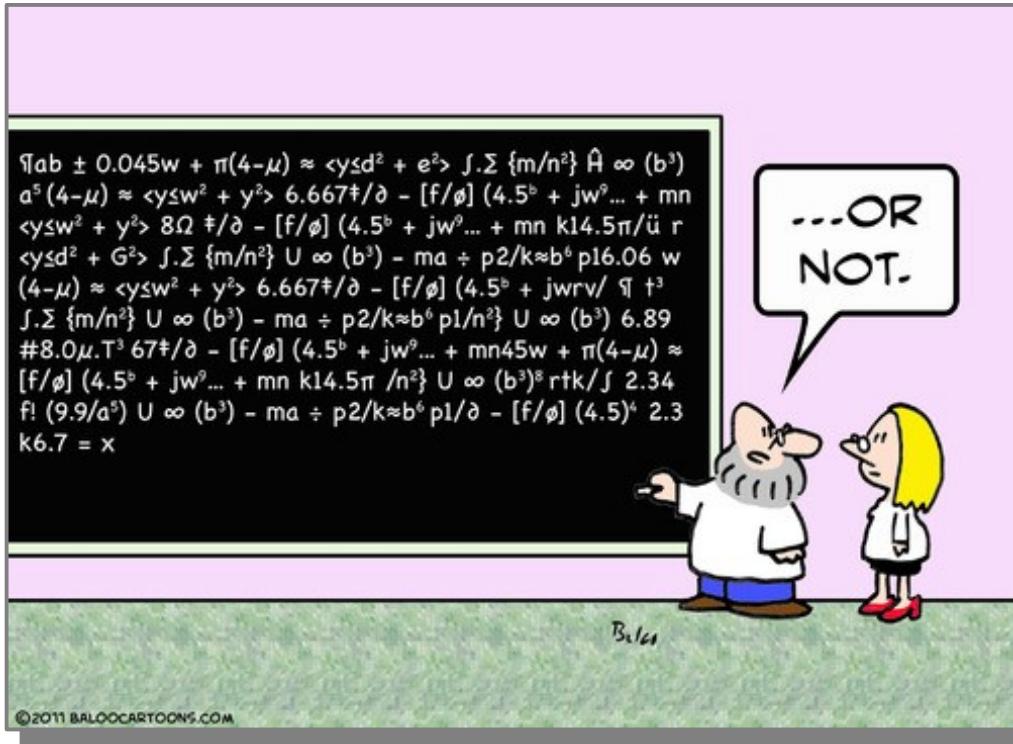


Double well potential  
on O-H bonds



Akrap et al, unpublished

# Tentative Theoretical Aspects



# Multiferroicity

6694

*J. Phys. Chem. B* 2000, 104, 6694–6709

## FEATURE ARTICLE

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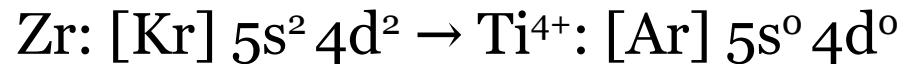
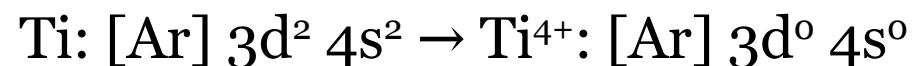
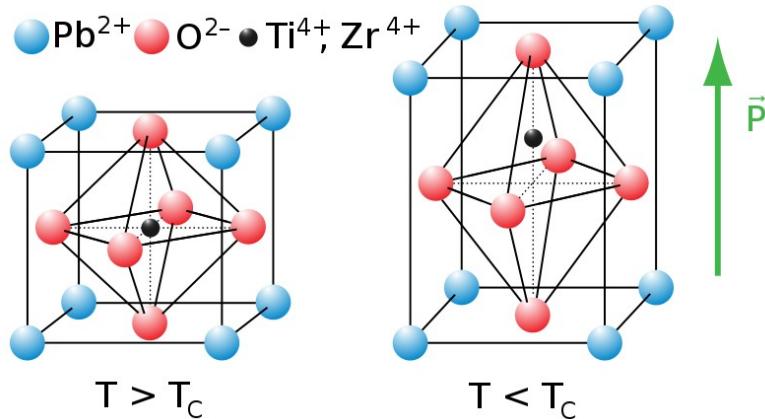
### Why Are There so Few Magnetic Ferroelectrics?

Nicola A. Hill

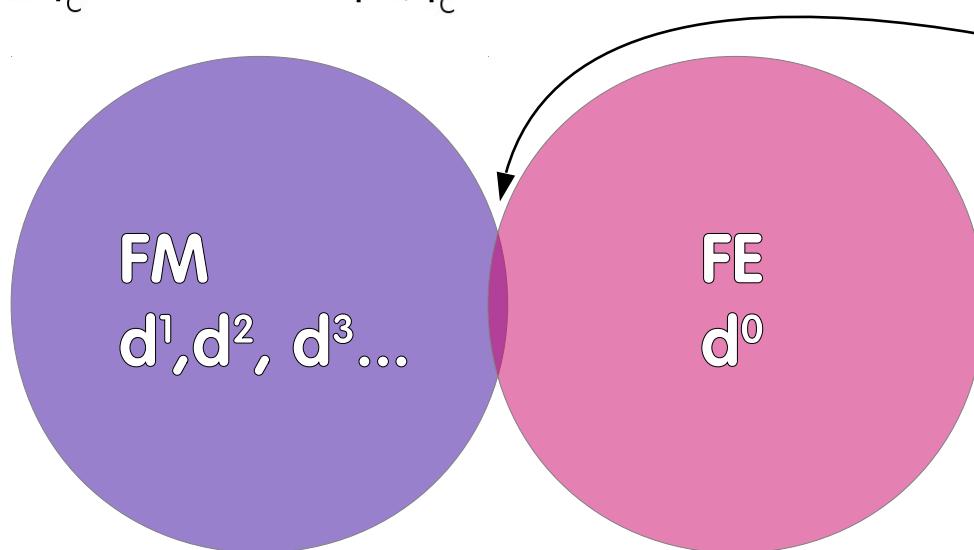
*Materials Department, University of California, Santa Barbara, California 93106-5050*

- ✓ Symmetry: 13 (out of 122) point groups allow spontaneous P and M to coexist
- ✓ Ferromagnetic materials like to be metallic (so AFM is welcome to the club)
- ✓ d<sup>o</sup>-ness

# $d^0$ -ness and Ferroelectricity



Ferroelectric distortions in perovskites like  $d^0$  electronic configurations



Multiferroics

Ab initio and observation shows that filling  $d$  orbitals inhibits FE distortions in perovskites

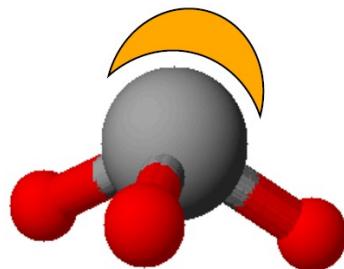
N. Hill (Spaldin), J. Phys. Chem. B **104**, 6694 (2000)

# Lone Pairs Ferroelectricity

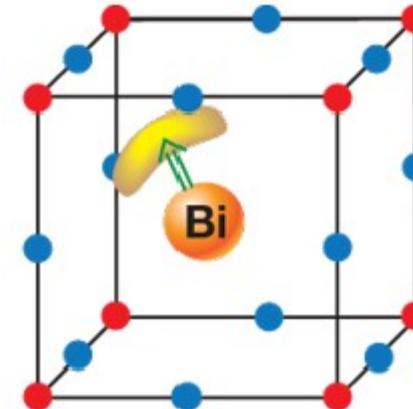
Pb: [Xe] 4f<sup>14</sup> 5d<sup>10</sup> 6s<sup>2</sup> 6p<sup>2</sup>

Bi: [Xe] 4f<sup>14</sup> 5d<sup>10</sup> 6s<sup>2</sup> 6p<sup>3</sup>

s<sup>2</sup> → “Dangling bonds”



p orbitals handle the chemical bonding



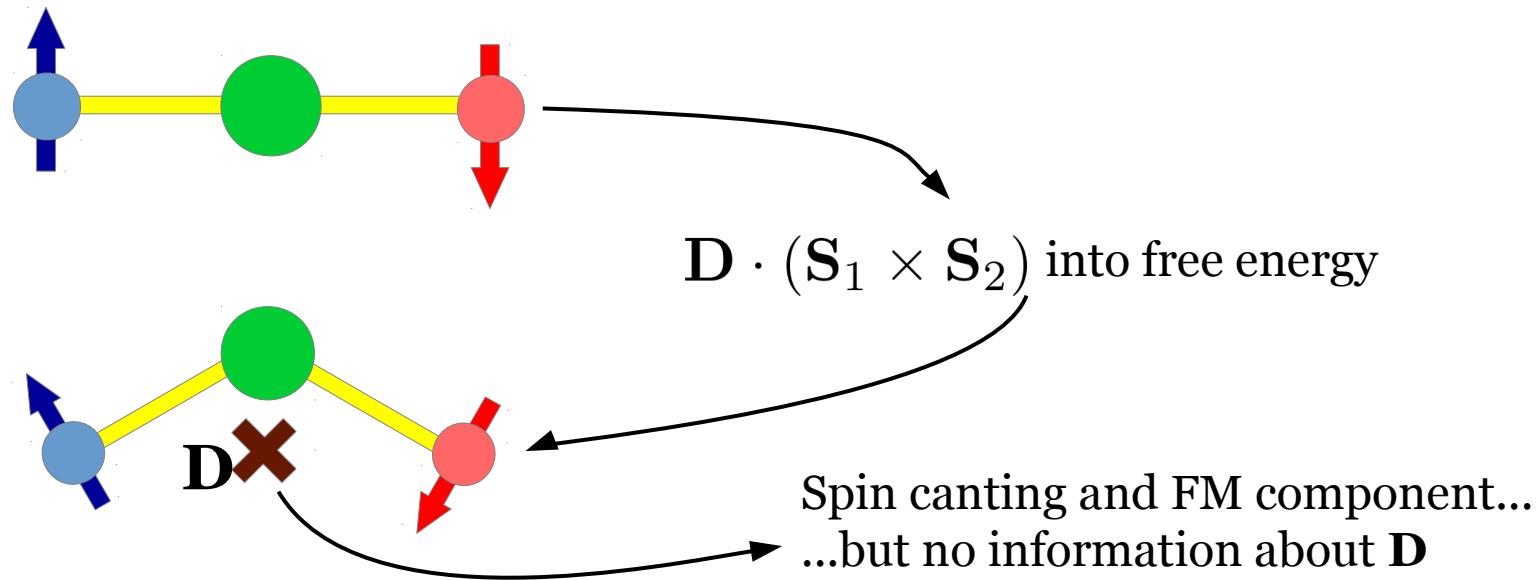
- High polarizability of s<sup>2</sup> electrons
- Local dipoles
- Main mechanism in BiFeO<sub>3</sub>
- Enhances ferroelectricity in PZT

BiFeO<sub>3</sub> beats d<sup>o</sup>-ness:

Ferroelectricity in A site (Bi)  
Antiferromagnetist in B site (Fe)

# (Inverse) Dzyaloshinskii-Moriya

Dzyaloshinskii: Why weak ferromagnetism in AFM materials?



Moriya:

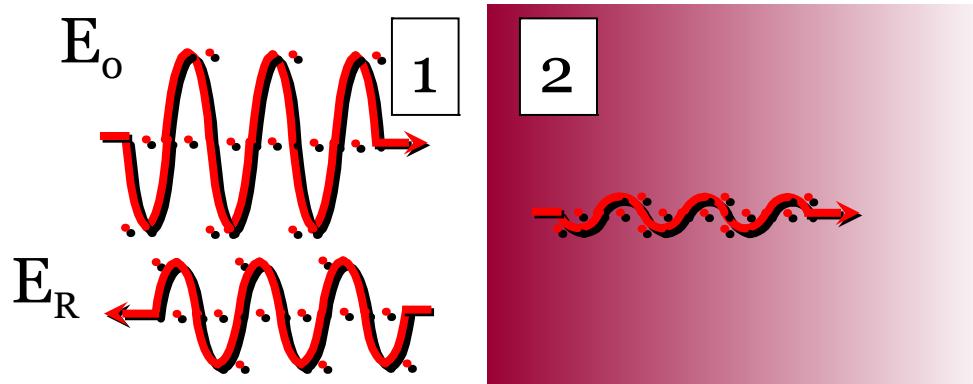
A diagram showing three spins (blue, green, red) connected by yellow rods. The green spin is at the top, the blue spin is at the bottom-left, and the red spin is at the bottom-right. A vector  $\mathbf{x}$  points upwards from the green spin, and a vector  $\mathbf{r}$  points from the blue spin to the red spin. A curved arrow labeled  $\mathbf{D} \propto \mathbf{r} \times \mathbf{x}$  points from this diagram to the text "Inverse DM: If the spins are already canted...".

**Inverse DM:** If the spins are already canted...

# Spectroscopical Toolbox



# Infrared



$\Delta \vec{k} \approx 0$  Momentum averaged

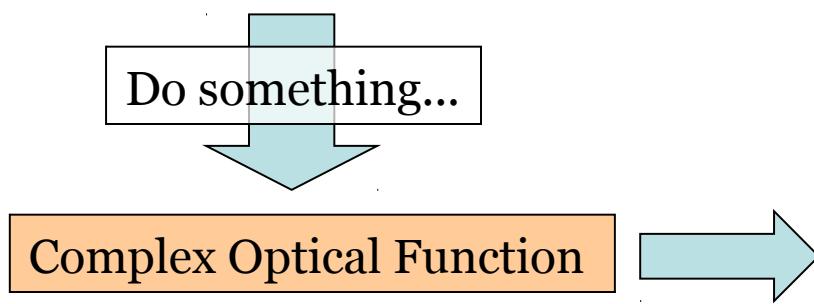
$\omega = \omega'$  Elastic scattering

$\omega \in [0, \infty]$  Broadband spectroscopy

$\chi_{Elec} \gg \chi_{Mag}$  Electric field dominates

$$R = \left| \frac{E_R}{E_0} \right|^2 = \left| \frac{1 - \sqrt{\varepsilon}}{1 + \sqrt{\varepsilon}} \right|^2 = \left| \frac{1 - n}{1 + n} \right|^2$$

Reflected Power (Real)



$$\begin{aligned} \sigma &= \sigma_1 + i\sigma_2 \\ &= i\varepsilon_0\omega(1 - \varepsilon) \end{aligned}$$

# Raman explained by Infrared people

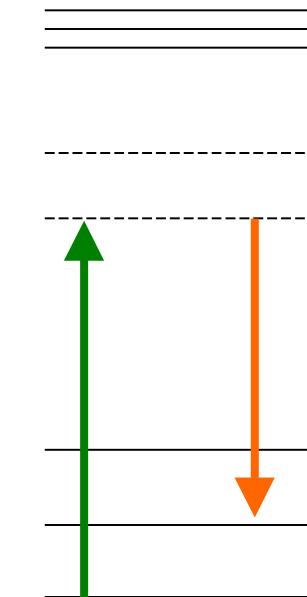
Excited elec. states

Virtual states

Excitation



Vibrational states



Anti-Stokes

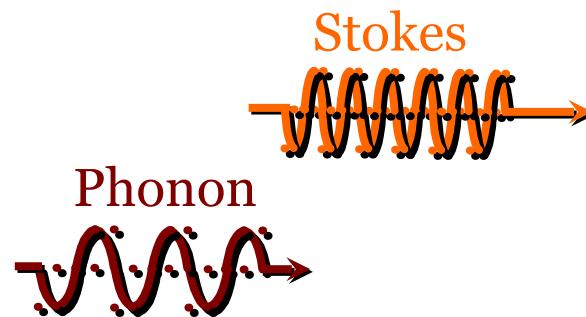
Rayleigh

Photon Energy

Excitation  
Energy

Stokes

Raman shift



# Infrared vs. Raman

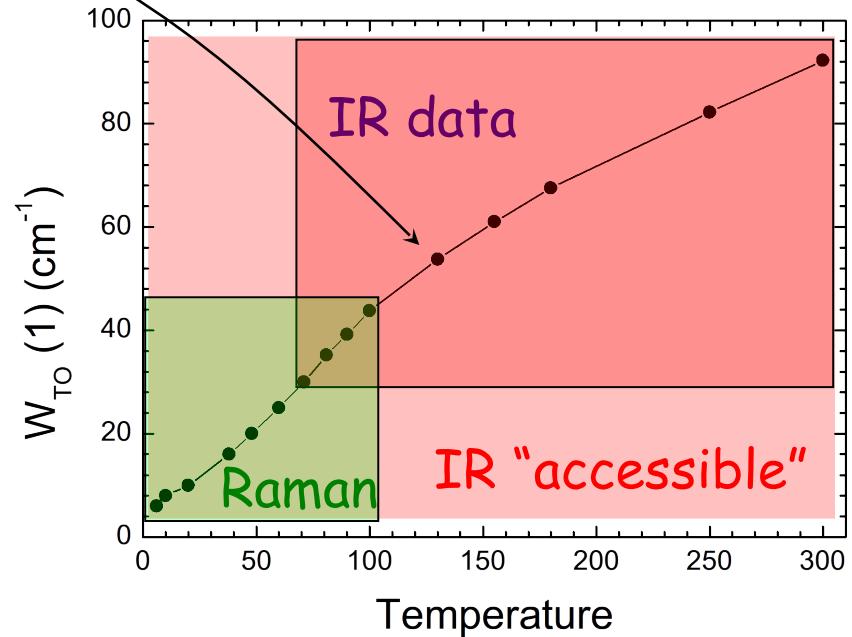
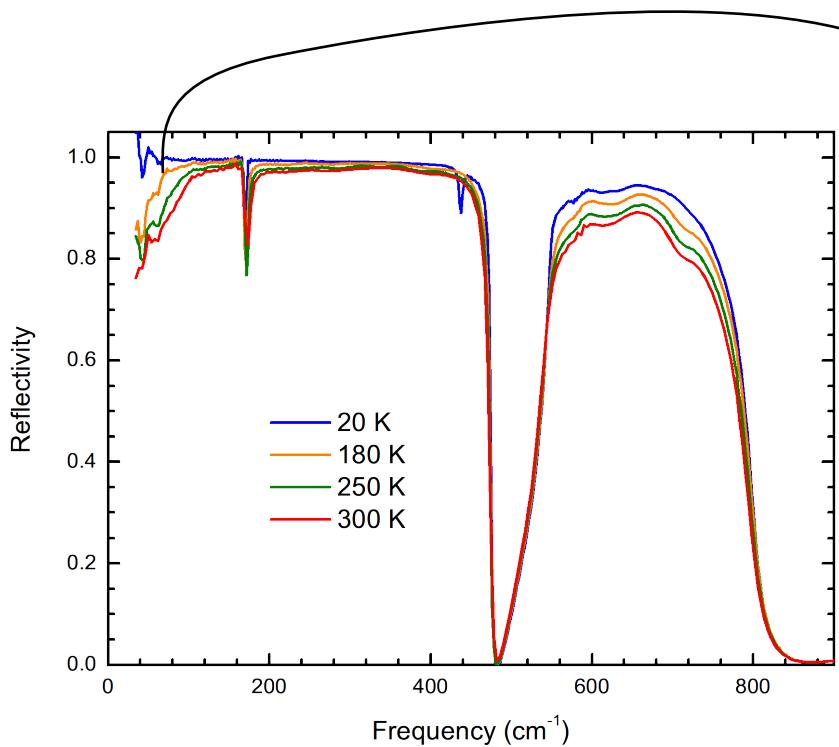
$$\mu = \alpha E$$

$\mu$  – Dipole moment → Infrared

$\alpha$  – Polarizability (2<sup>nd</sup> order tensor) → Raman

32 Point Groups	Classes	IR vs. Raman
$C_1, C_2, C_s, C_3, C_{3v}$	5 polars (= pyro)	Raman = IR
$C_{2v}, C_4, C_{4v}, C_6, C_{6v}$	5 polars (= pyro)	Raman ⊃ IR
$D_2, S_4, D_{2d}, T, T_d$	5 piezo, non pyro	Raman ⊃ IR
$D_4, D_3, C_{3h}, D_6, D_{3h}$	5 piezo, non pyro	E in Raman & IR; A in IR; Other in Raman
$C_{2h}, D_{2h}, C_{4h}, D_{4h}, S_6,$ $D_{3d}, C_{6h}, D_{6h}, T_h, O_h, O$	11 inversion symmetry and O	Raman ∩ IR = $\emptyset$

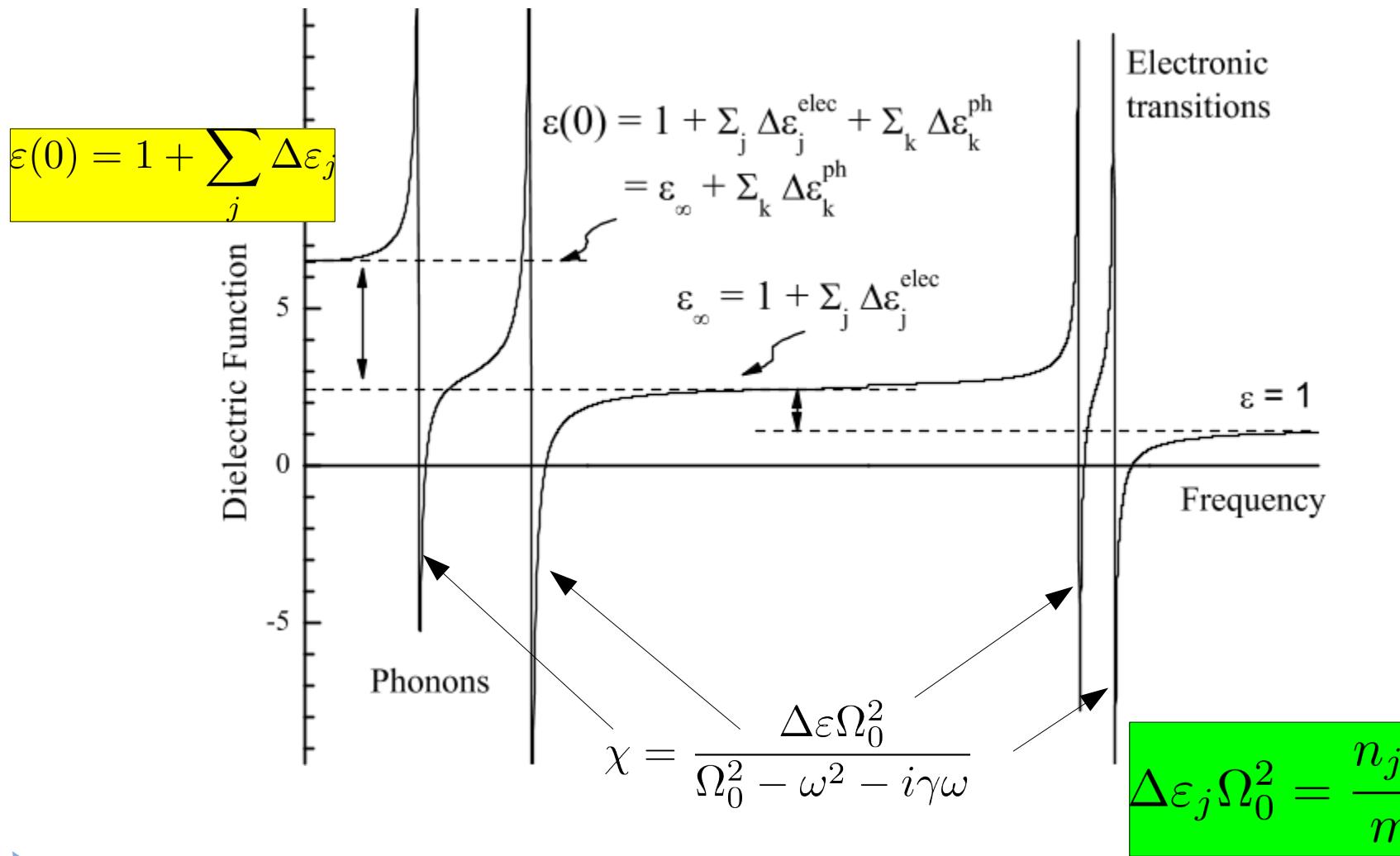
# $\text{SrTiO}_3$ – IR / Raman working together



**Bad:** Long wavelengths (low energies) in IR hit diffraction  
(technical only) problems.

**Good:** IR access the polar phonons in high T and low T symmetries

# Building the Dielectric Constant



# The f-sum rule

(Pour les voyageurs dans le temps, voir cours de Vihn samedi et revenir...)

The f-sum rule (particle conservation):

$$\int_0^\infty \sigma_1(\omega) d\omega \sim \int_0^\infty \omega \varepsilon''(\omega) d\omega \sim \frac{n}{m}$$

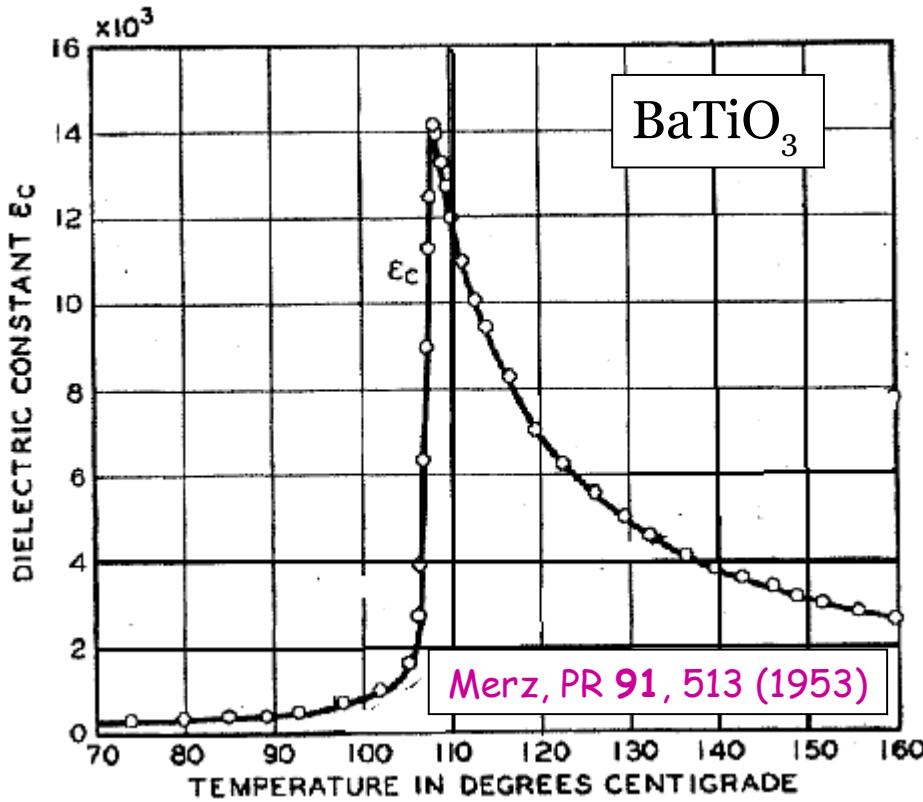
The f-sum rule for phonons:

$$\sum_j \Delta\varepsilon_j \Omega_{0j}^2 = \sum_j \frac{n_j q_j^2}{m_j} = const$$

The f-sum rule for decoupled phonons:

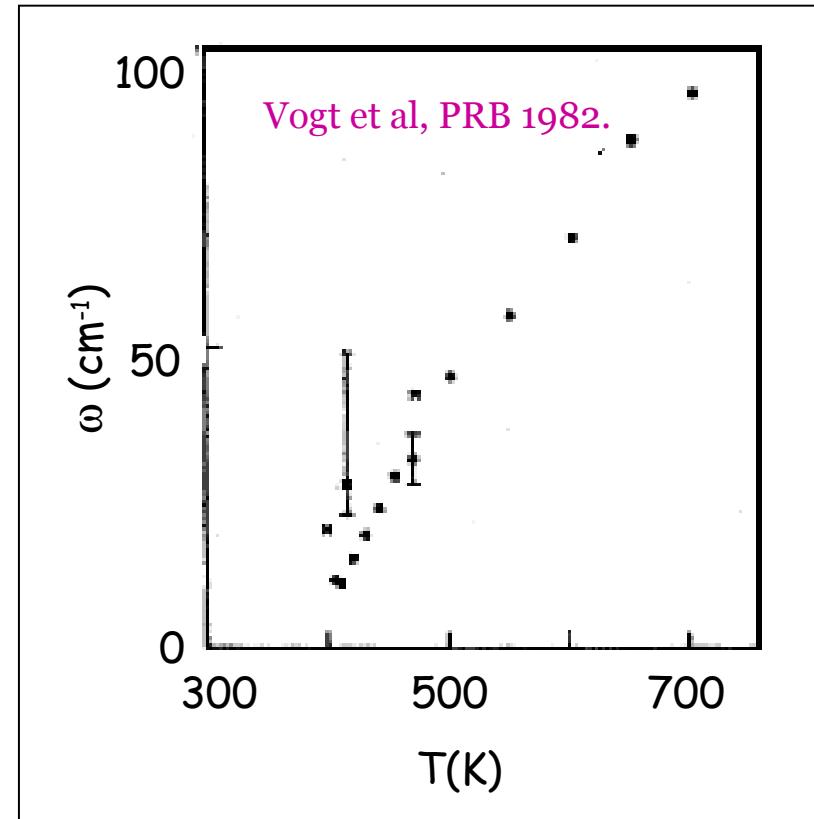
$$\Delta\varepsilon_j \Omega_{0j}^2 = \frac{n_j q_j^2}{m_j} = const$$

# Soft mode and Ferroelectrics

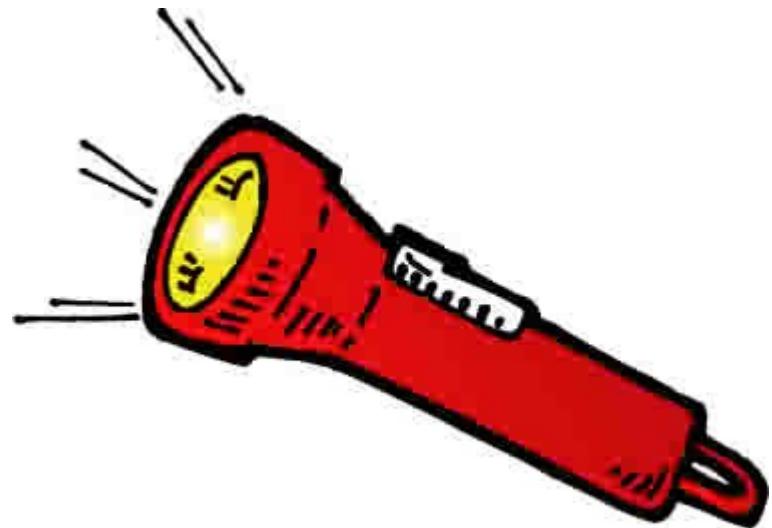


$$\epsilon(0) = 1 + \sum_j \Delta\epsilon_j$$

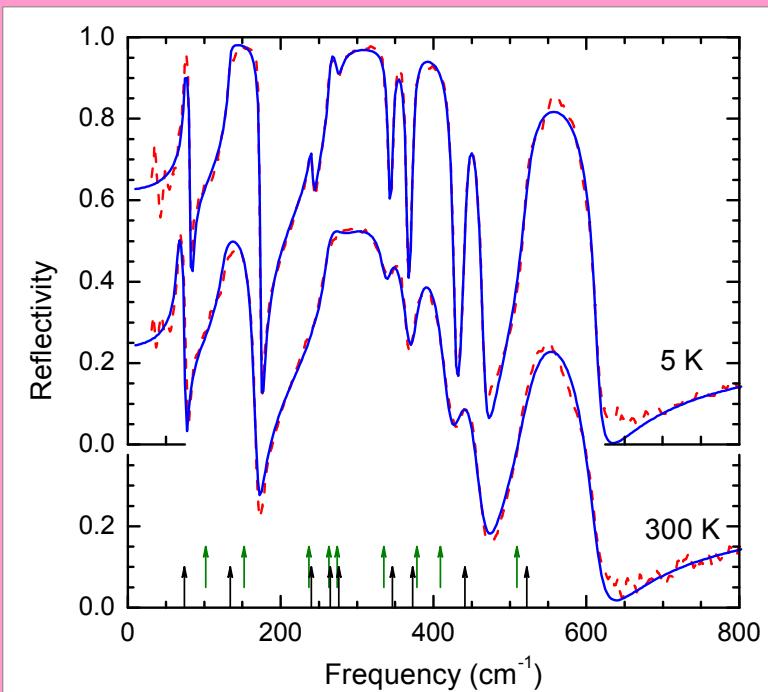
$$\Delta\epsilon_j \Omega_{0j}^2 = \frac{n_j q_j^2}{m_j} = \text{const}$$



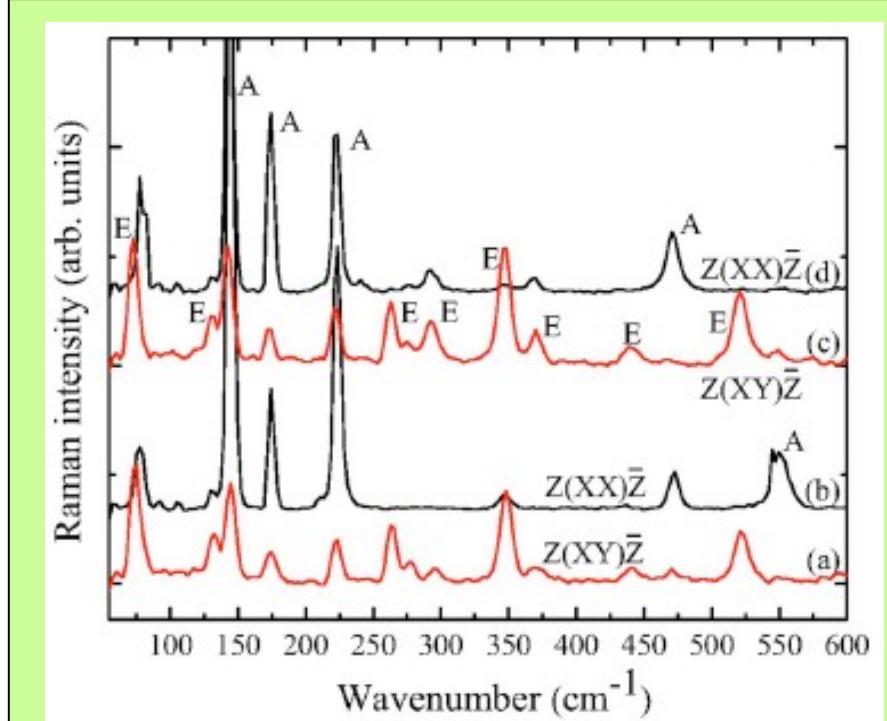
# Spectroscopy of Multiferroics



# Raman & IR on BiFeO<sub>3</sub>



Lobo et al. PRB **76**, 172105 (2007)

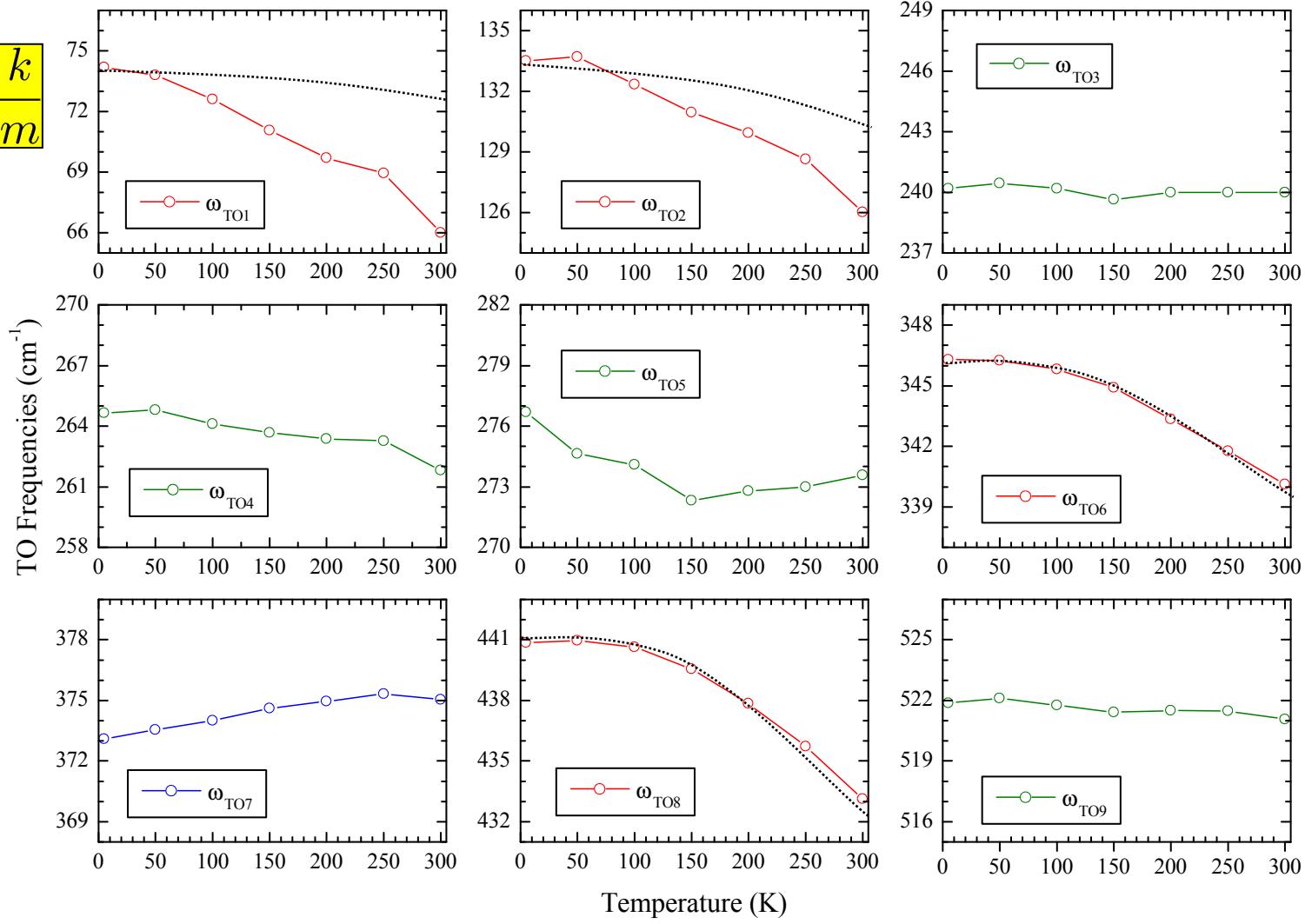


Cazayous et al. APL **91**, 071910 (2007)

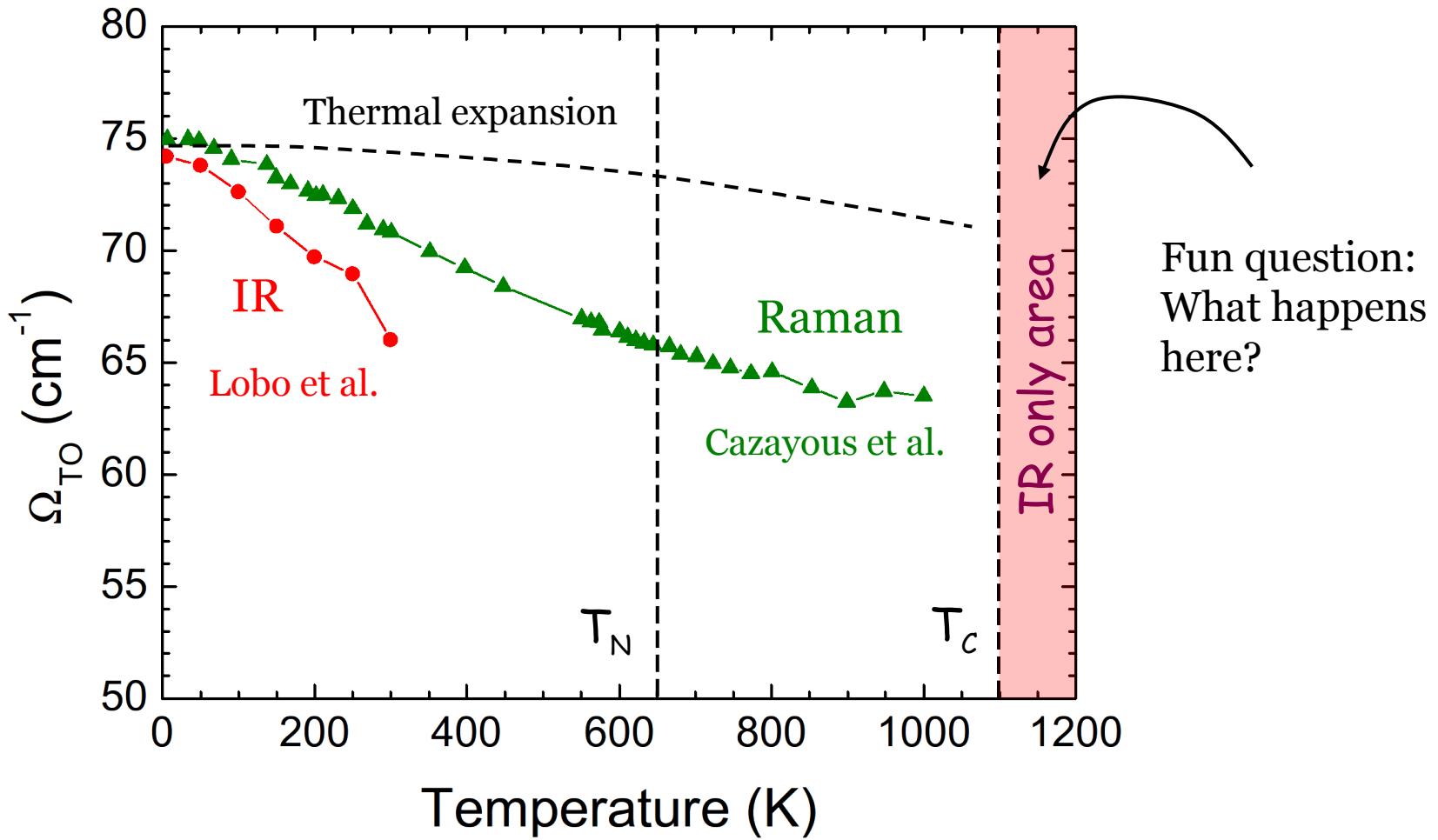
- ✓ Number of modes as predicted by group theory
- ✓ Good correlation with ab initio

# Transverse in-plane optical modes

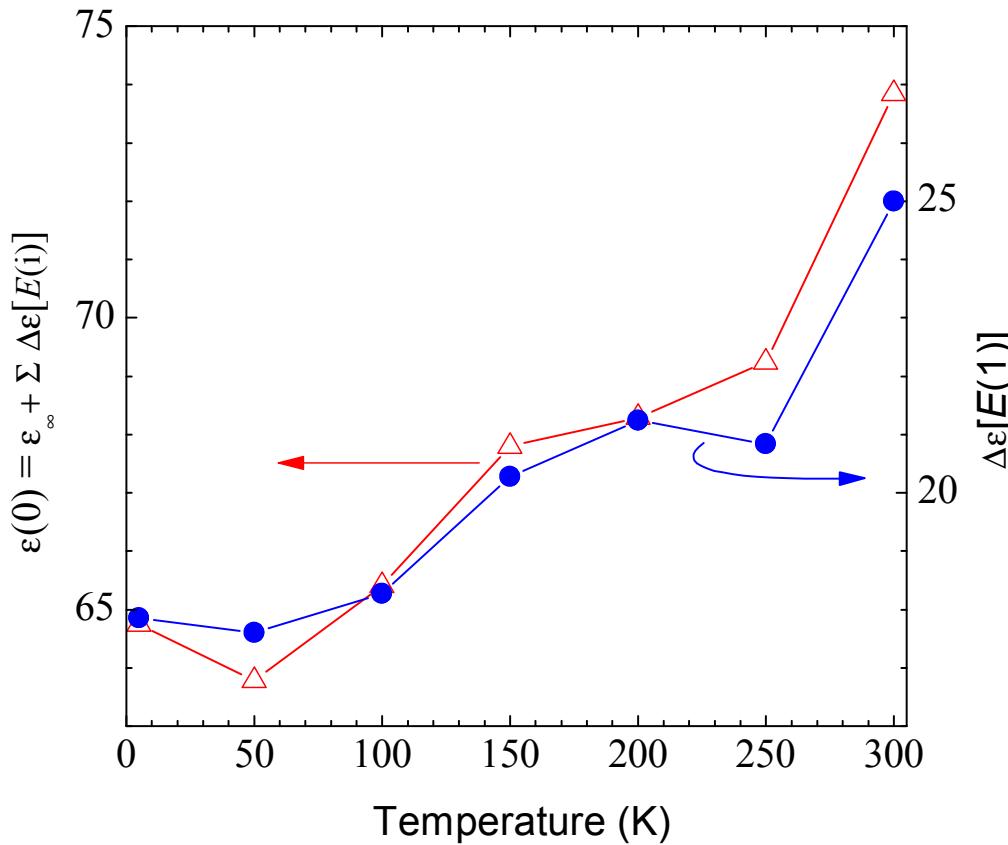
$$\Omega_0^2 \sim \frac{k}{m}$$



# Hunting the Soft Mode

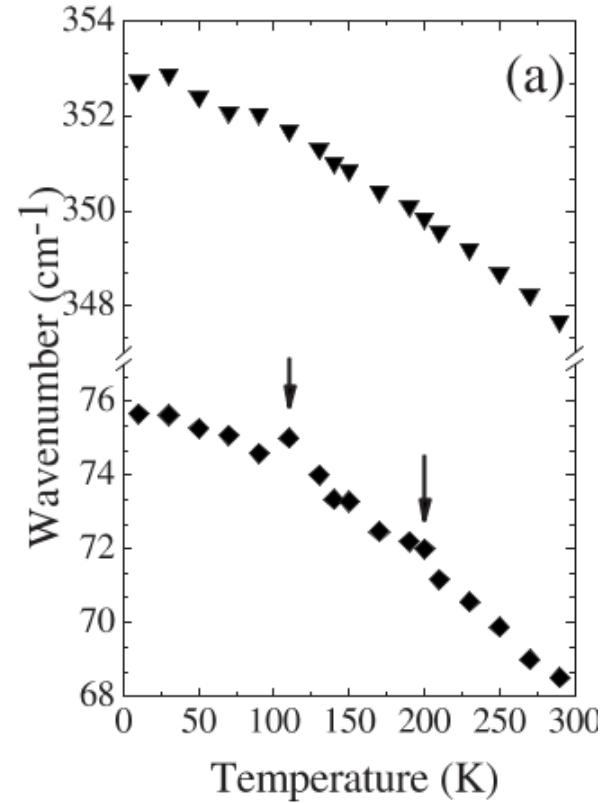
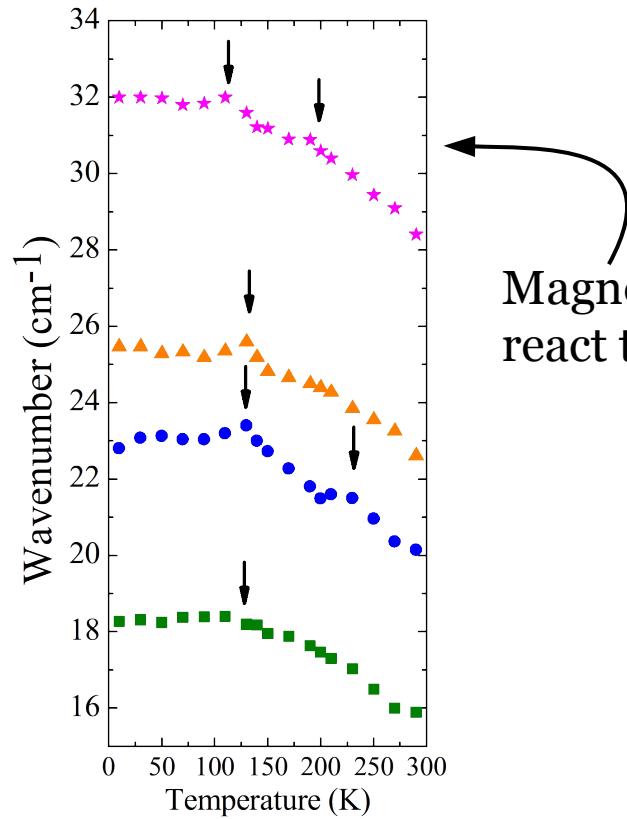


# The Soft Mode Controls the Dielectric Constant



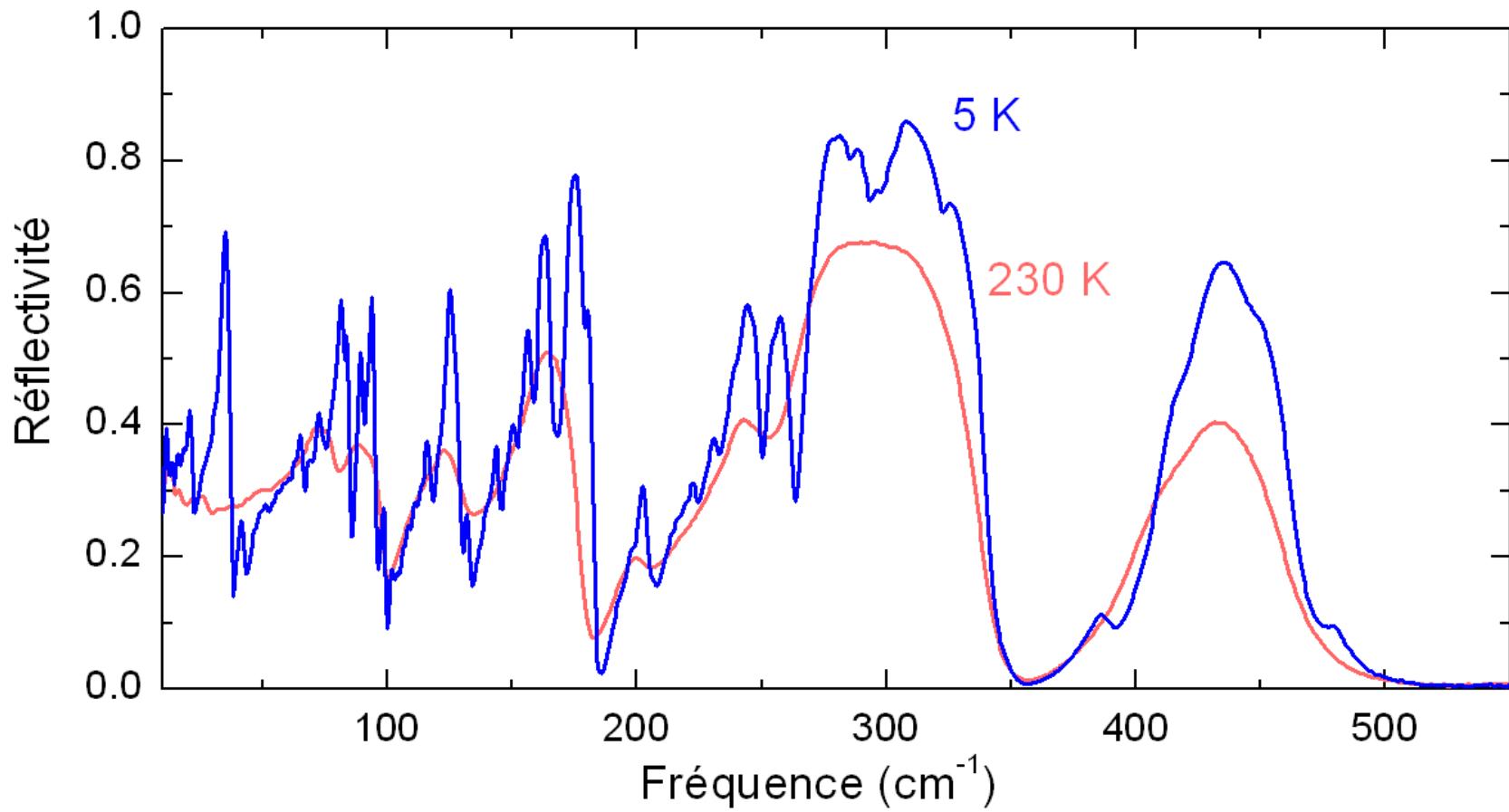
The temperature dependence of the static dielectric constant comes fully from the lowest  $E$  symmetry phonon.

# Dielectric (phono) Magnetic (magnon) Interaction



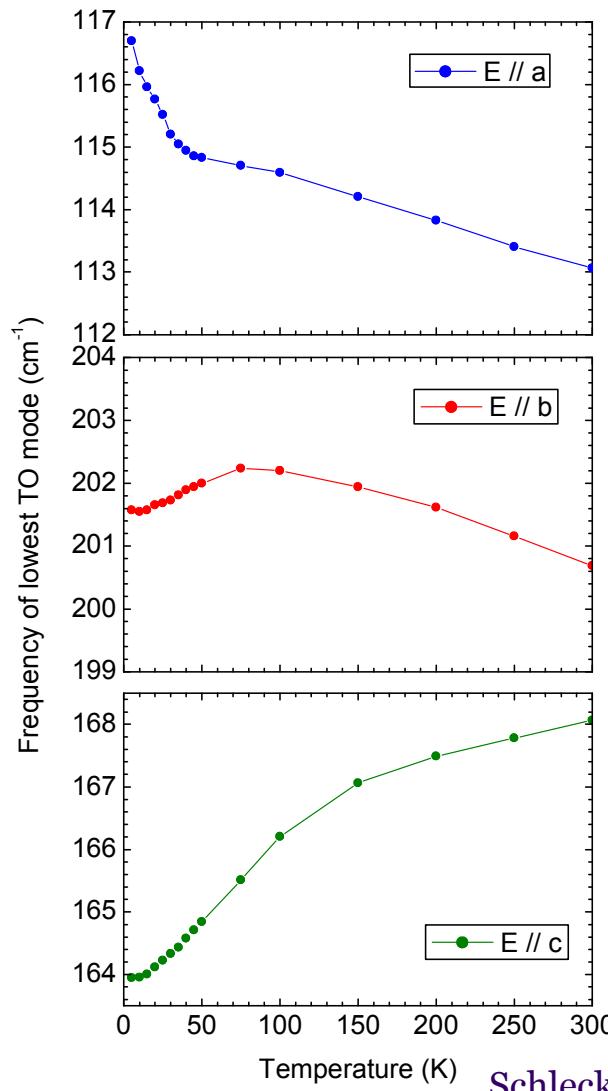
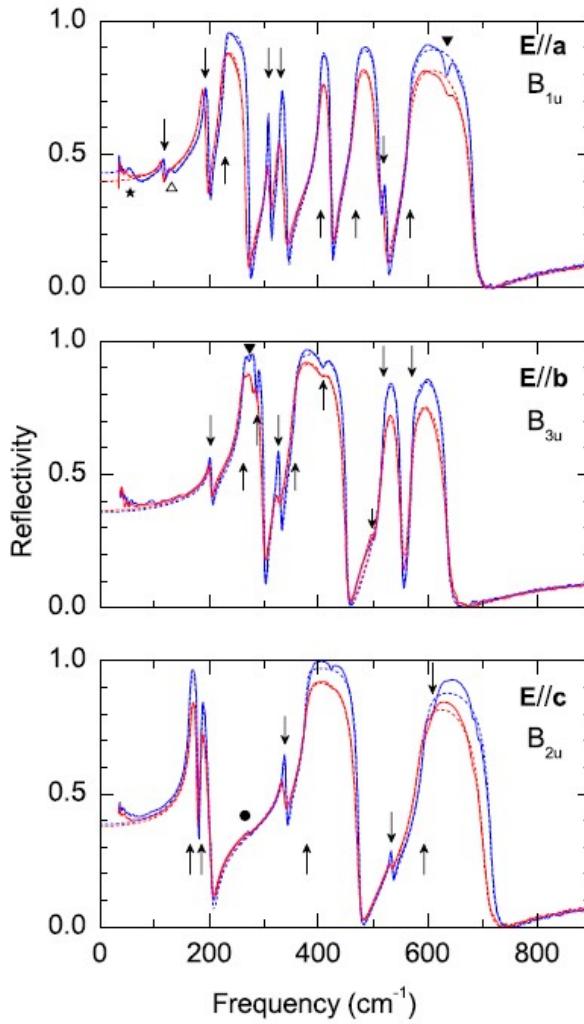
Rovillain et al., PRB **79**, 180411 (2009)

# Before looking into $\text{TbMnO}_3$ , remind $\text{BaMnF}_4$



Schleck et al. 2012.

# Infrared response of $\text{TbMnO}_3$



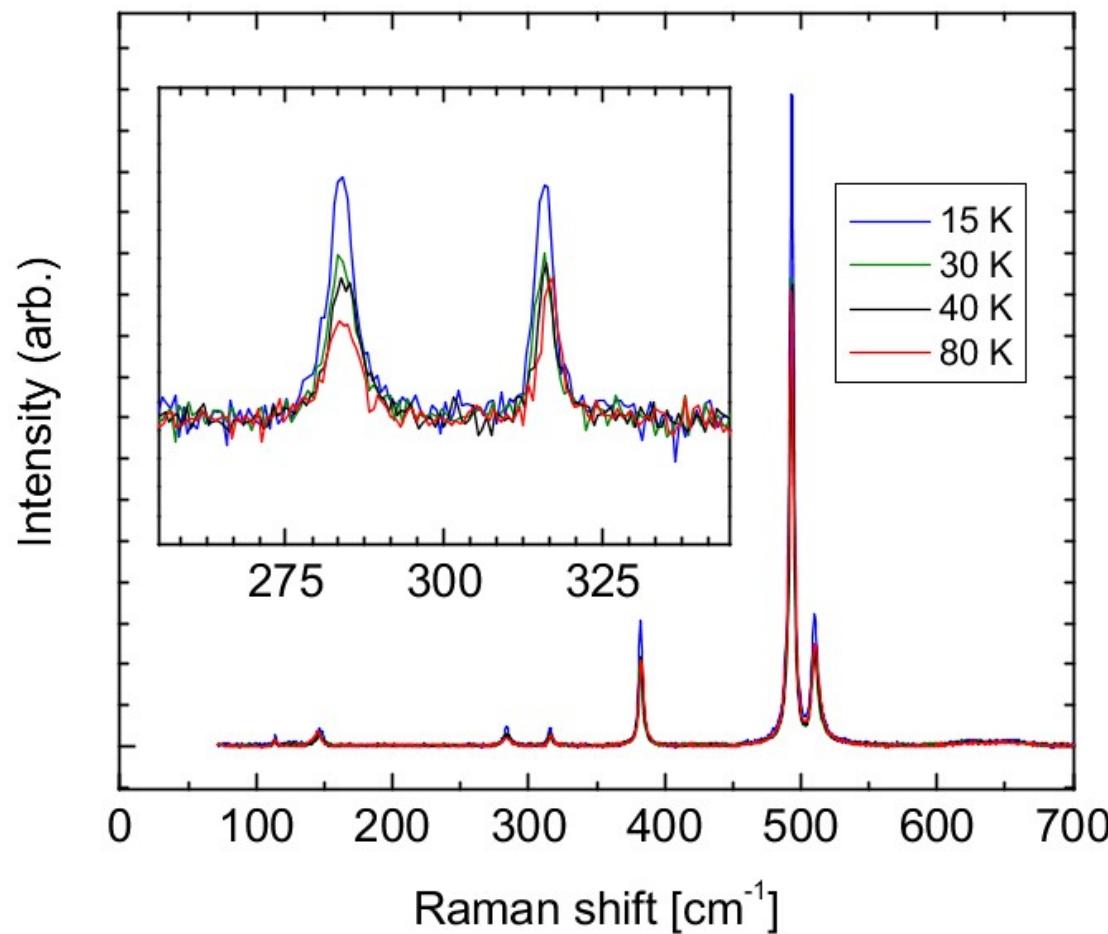
No sign of inversion symmetry breaking

No phonon spectra reconstruction

No soft mode

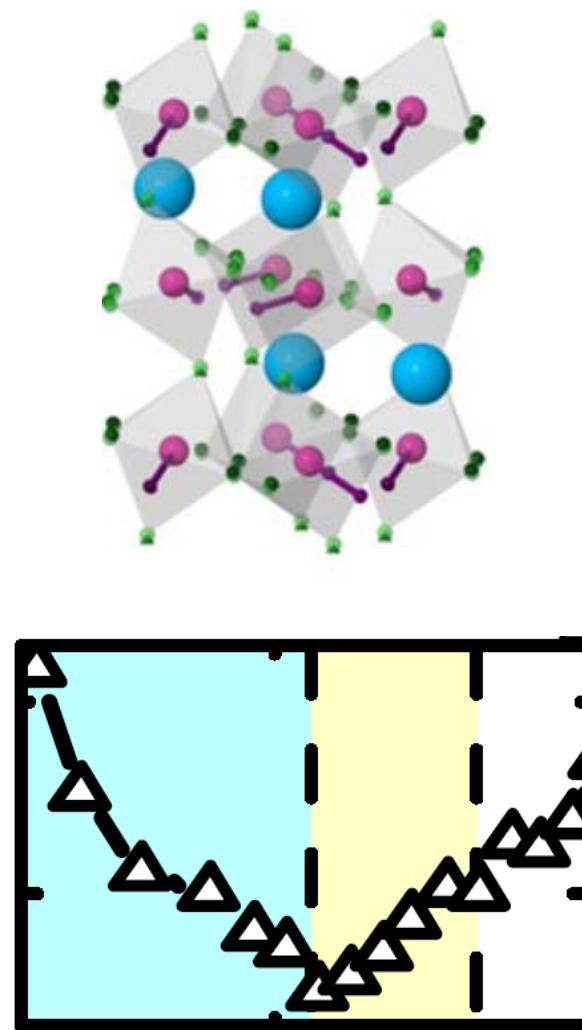
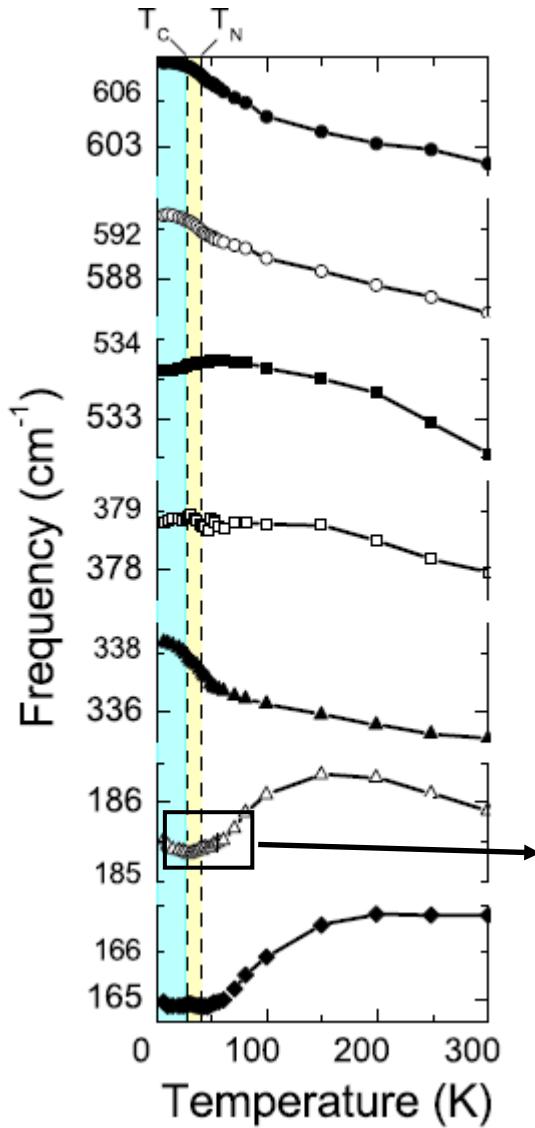
Schleck et al., PRB **82**, 144309 (2010)

# Raman in $\text{TbMnO}_3$ ain't any better



Rovillain, Cazayous, unpublished

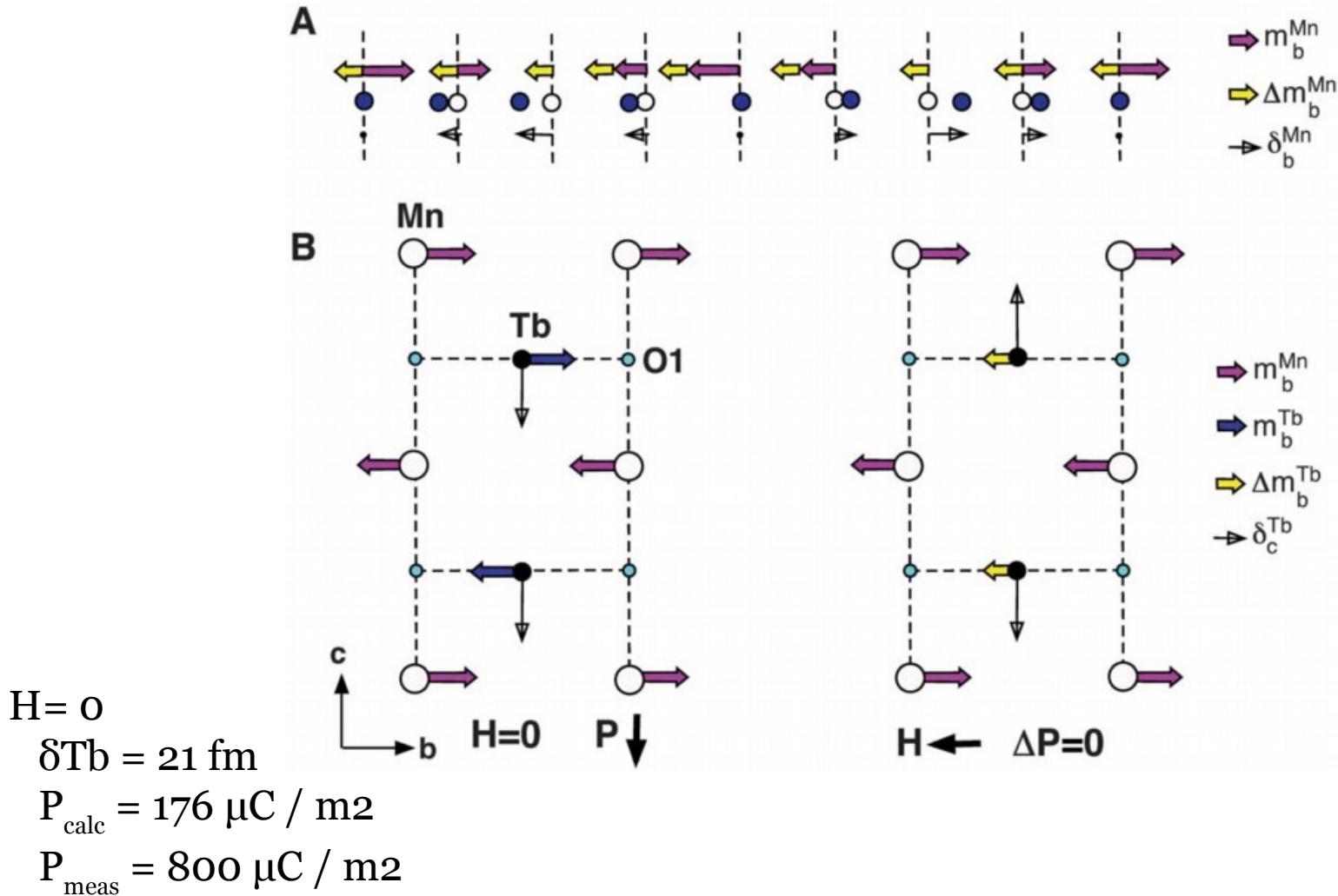
# Looking real close at $\text{TbMnO}_3$ phonons



- ✓ Phonons are slightly renormalized at the AF transition
- ✓ The only phonon that has (a very tiny) modification at the ferroelectric transition is dominated by Mn atomic motions.

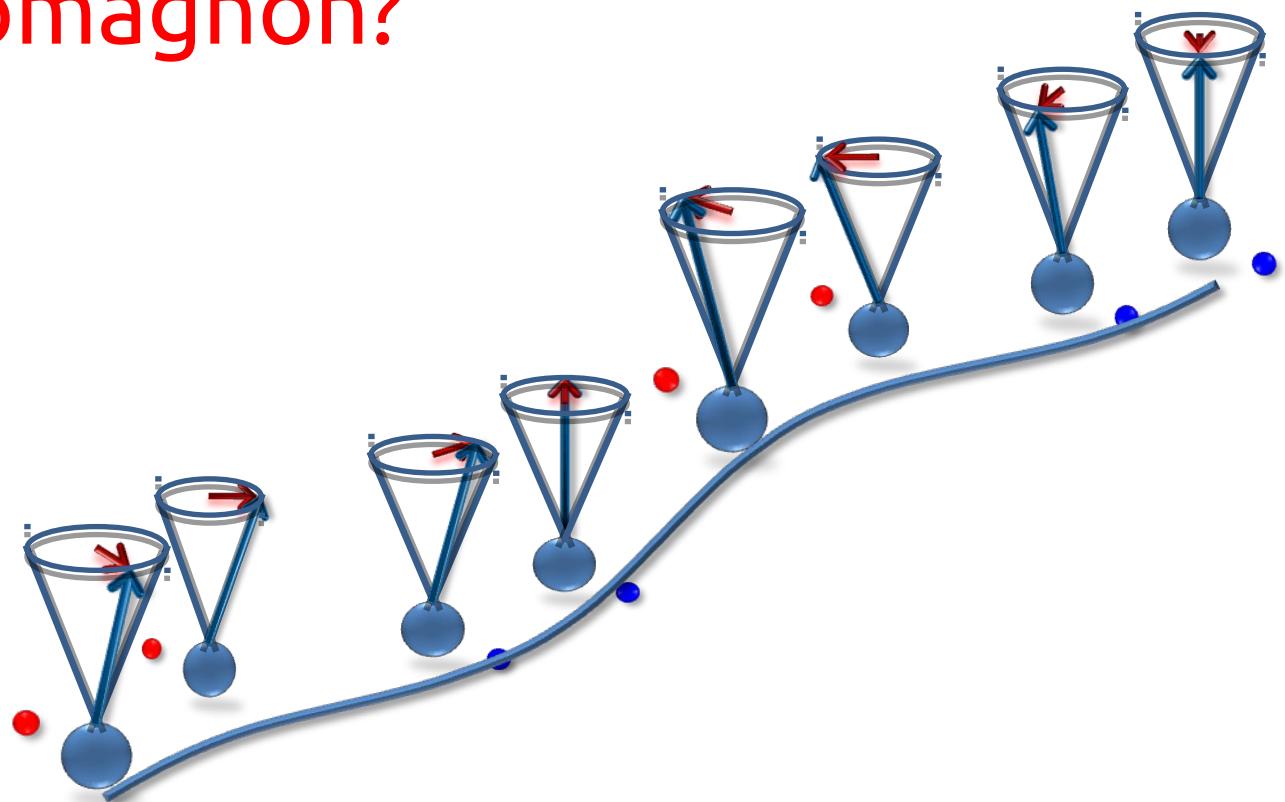
Schleck et al., PRB **82**, 144309 (2010)

# Is there an atomic displacement

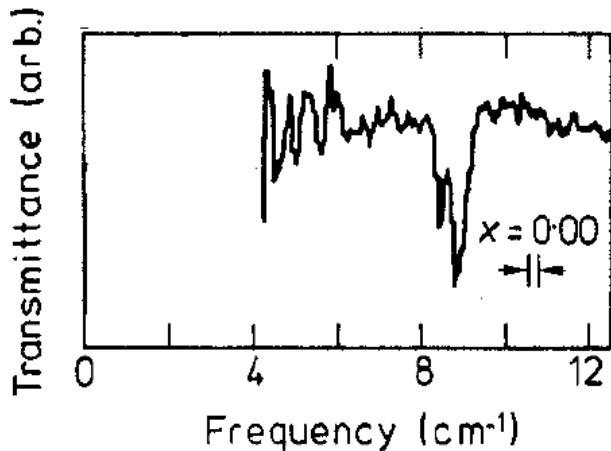


H. C. Walker et al. Science 333, 1273 (2011)

# What is an Electromagnon?

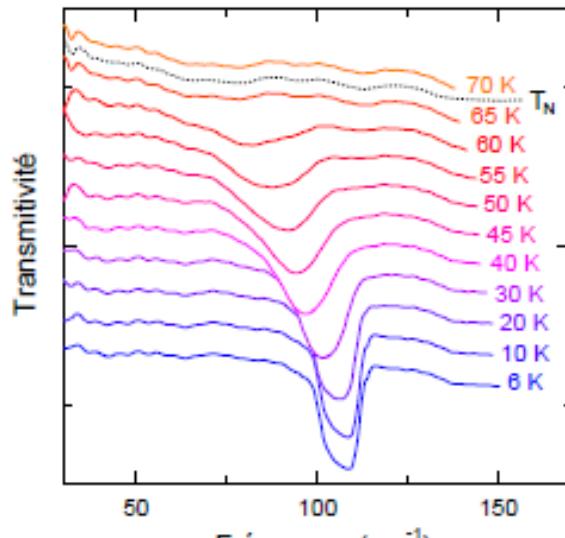


# Infrared Magnetic Excitations in $\text{MnF}_2$

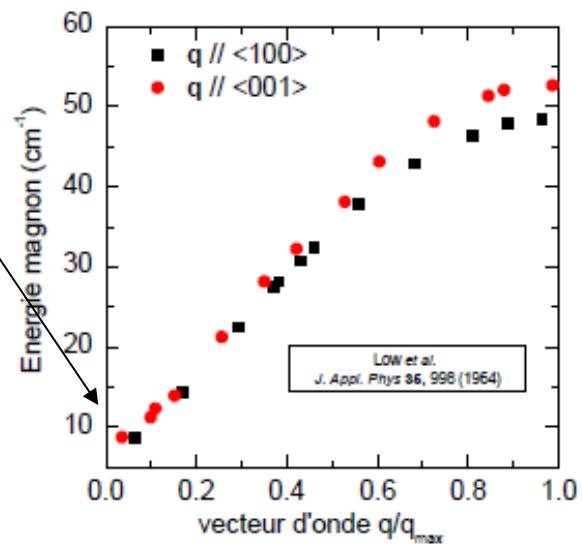


Allen et al., PRL 16, 463 (1966)

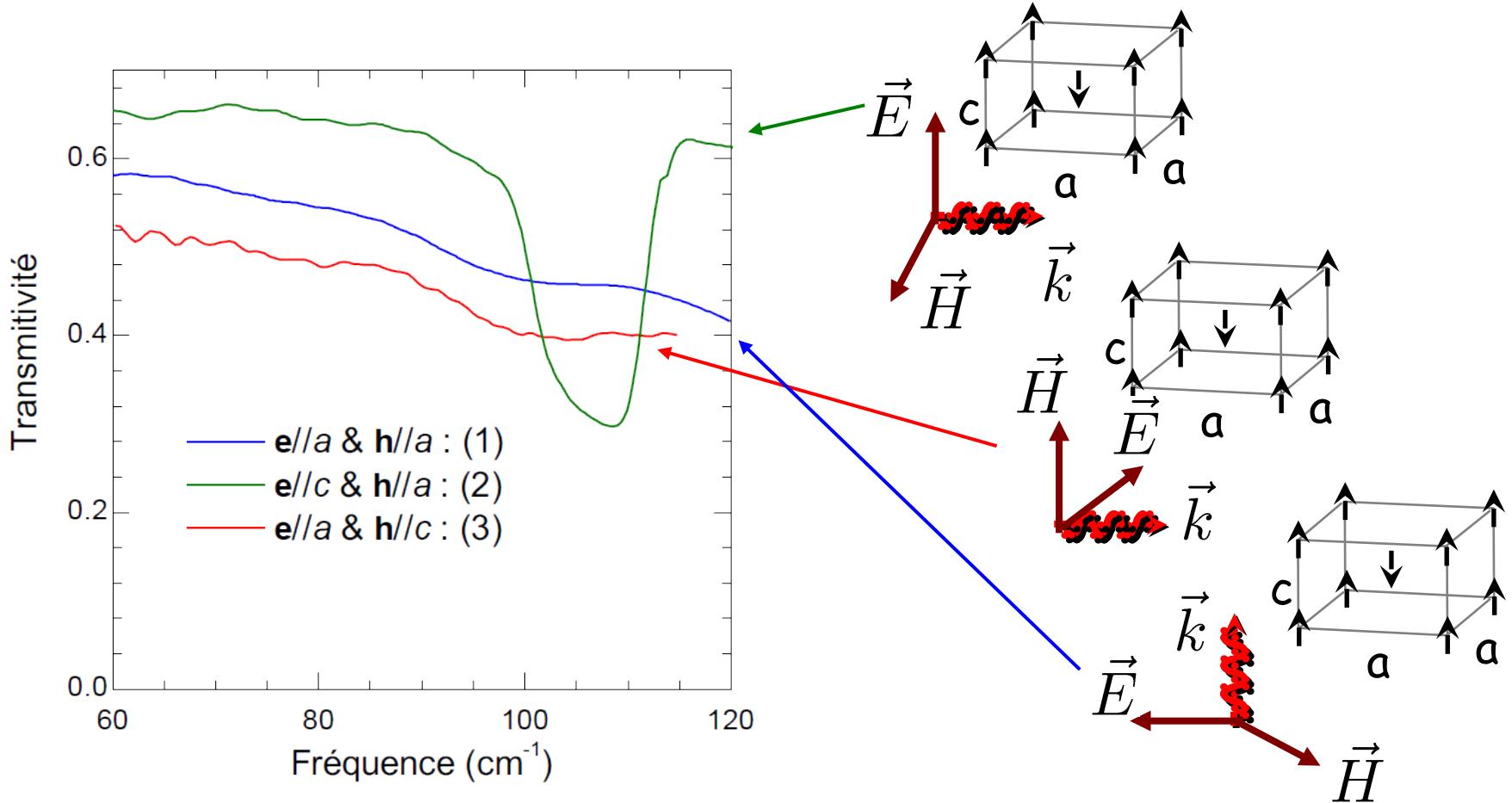
AFM resonance



Zone edge bi-magnon

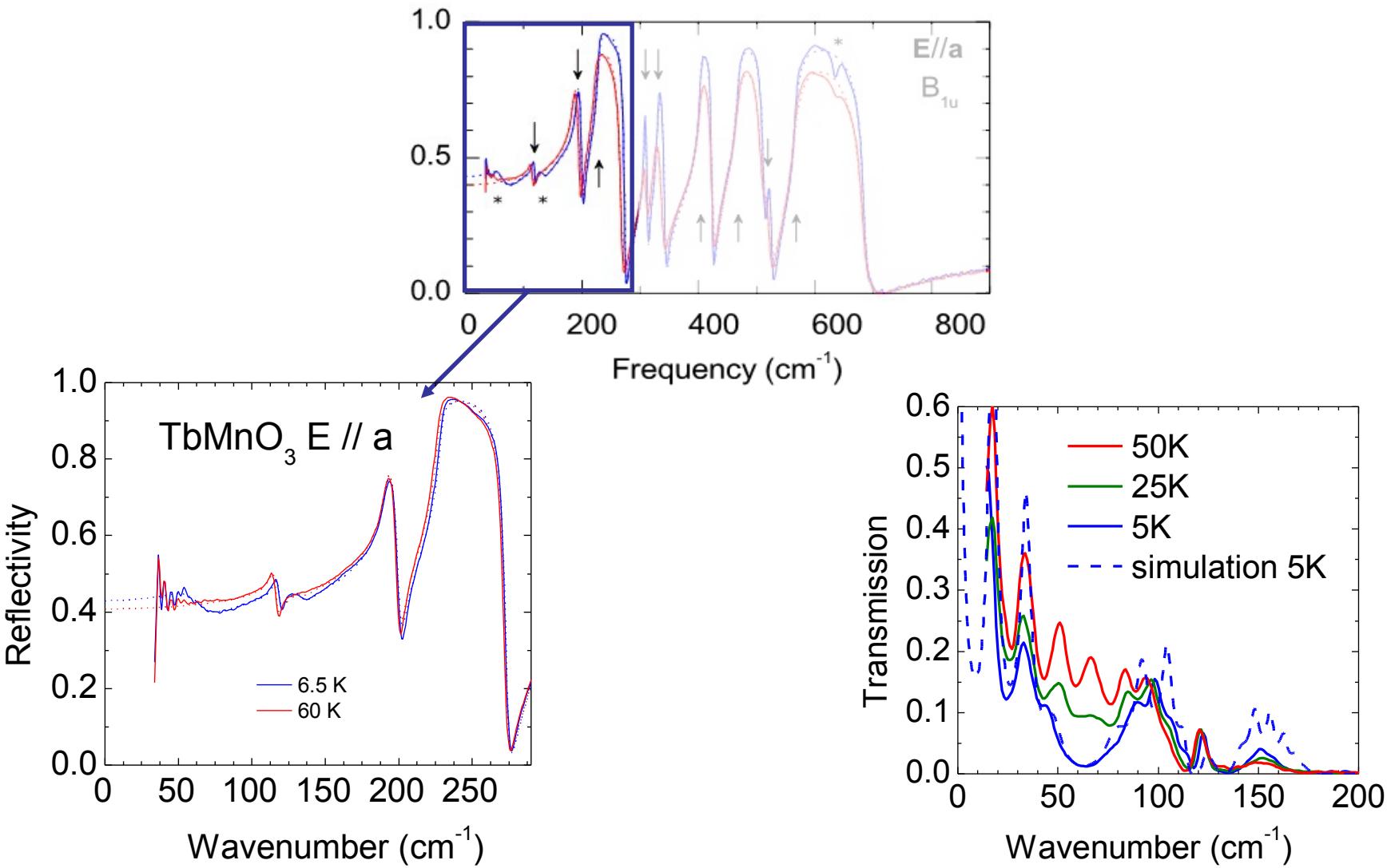


# How to determine if an excitation is electric



The bi-magnon in  $\text{MnF}_2$  is activated by the electric field of light.  
Not its magnetic field.

# Taking a closer look at $\text{TbMnO}_3$



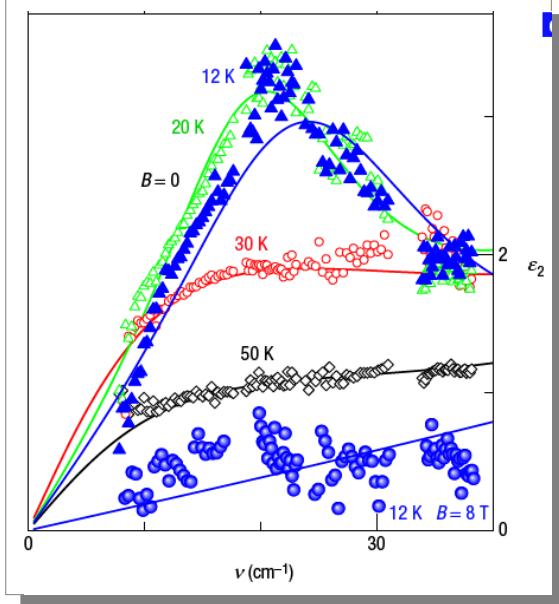
# Rough definition of the Electromagnon

Possible evidence for electromagnons in  
multiferroic manganites

Nat. Phys, 2006.

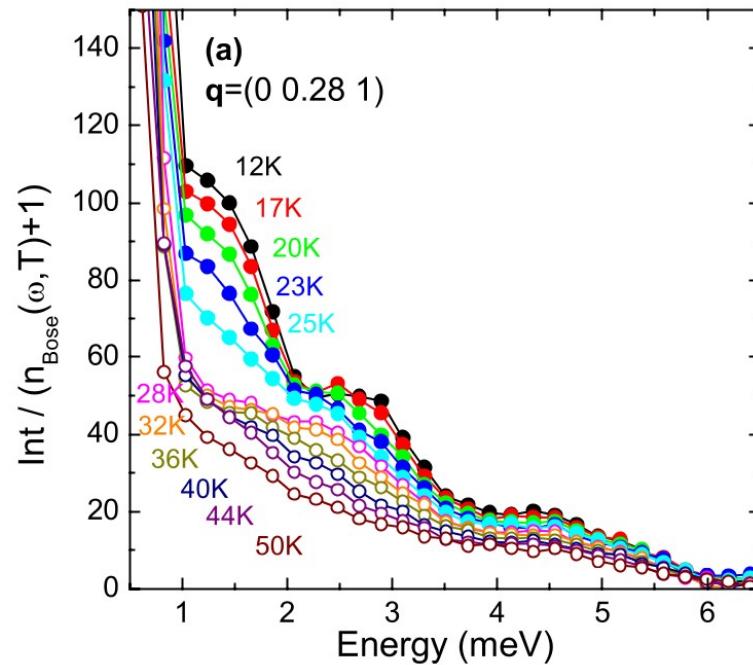
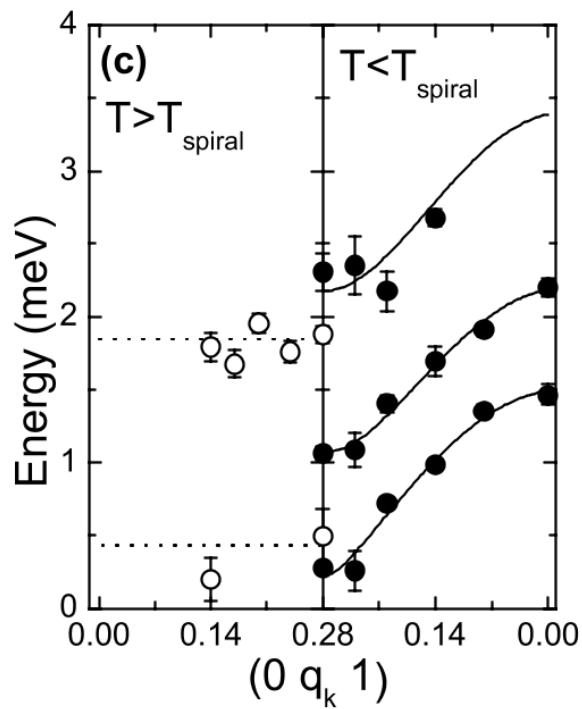
A. PIMENOV<sup>1\*</sup>, A. A. MUKHIN<sup>1,2</sup>, V. YU. IVANOV<sup>2</sup>, V. D. TRAVKIN<sup>2</sup>, A. M. BALBASHOV<sup>3</sup> AND A. LOIDL<sup>1</sup>

and polar order. Among other multiferroics,  $\text{TbMnO}_3$  and  $\text{GdMnO}_3$  reveal a strong magneto-dielectric coupling and as a consequence fundamentally different spin excitations exist: electro-active magnons (or electromagnons), spin waves that can be excited by a.c. electric fields. Here we provide evidence that



- ✓ Spin wave activated by electric field of light
- ✓ Is the electromagnon a new fundamental excitation?
- ✓ What is the relationship between the electromagnon and the magneto-electric coupling?
- ✓ Is the electromagnon a necessary excitation for ME coupling?
- ✓ Is the electromagnon restricted to multiferroics?

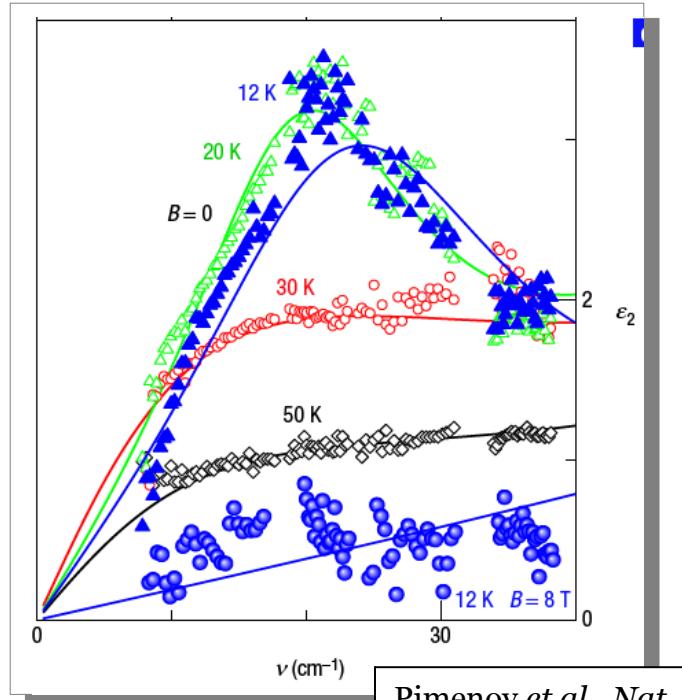
# The Electromagnon is also a Magnon



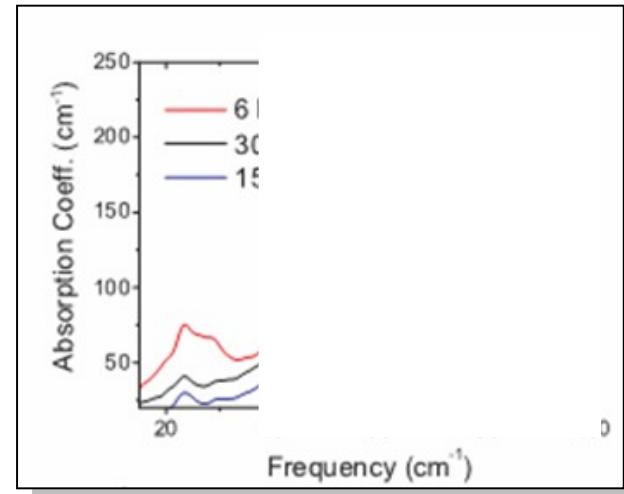
- Same energy of IR electromagnon
- Stronger below  $T_c$
- Persists up to  $T_N$

Senff et al., PRL **98**, 137206 (2007)

# Not just one e-magnon in $\text{TbMnO}_3$



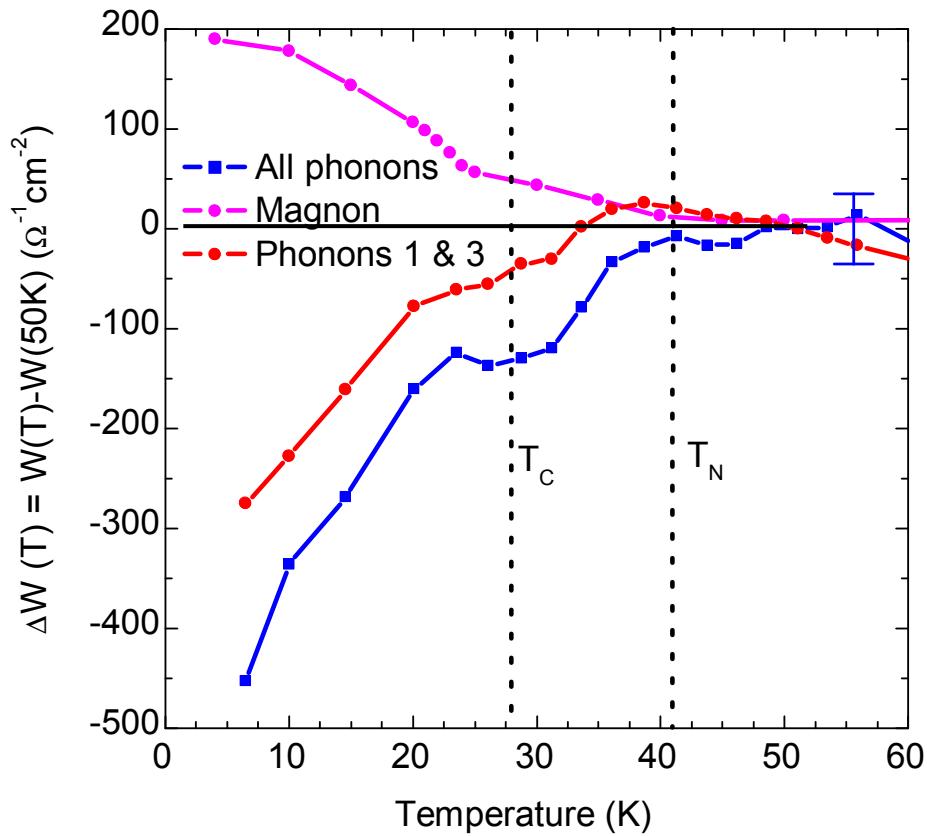
Pimenov *et al.*, *Nat. Phys.*, 2006.



Valdés Aguilar *et al.*, *PRL*, 2009.

Most of the e-magnon spectral weight is  
at higher frequencies

# The electromagnon ( $E//a$ ) is built from two phonons

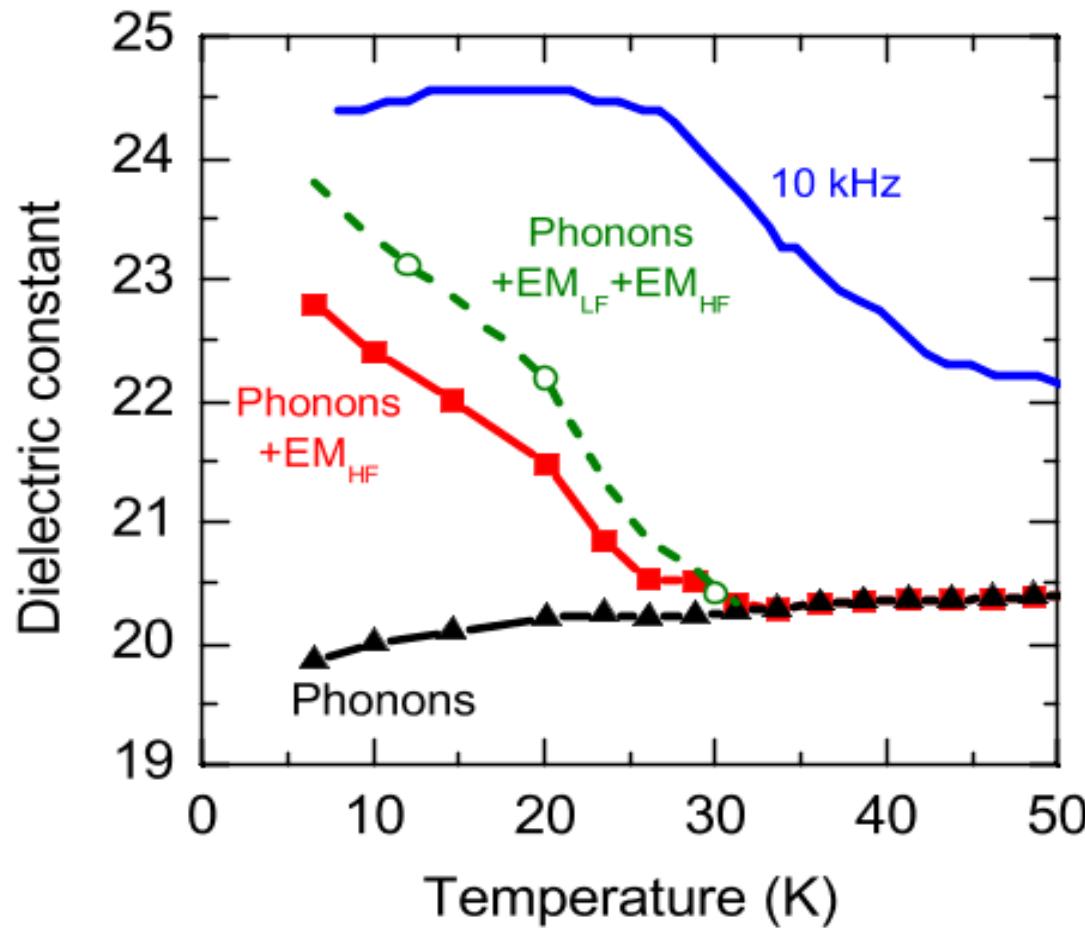


Sum-rule

$$\sum_j \Delta\varepsilon_j \Omega_{0j}^2 = \text{const}$$

The electromagnon in  $\text{TbMnO}_3$  gets its strength from phonons dominated by Tb and Mn atomic movements

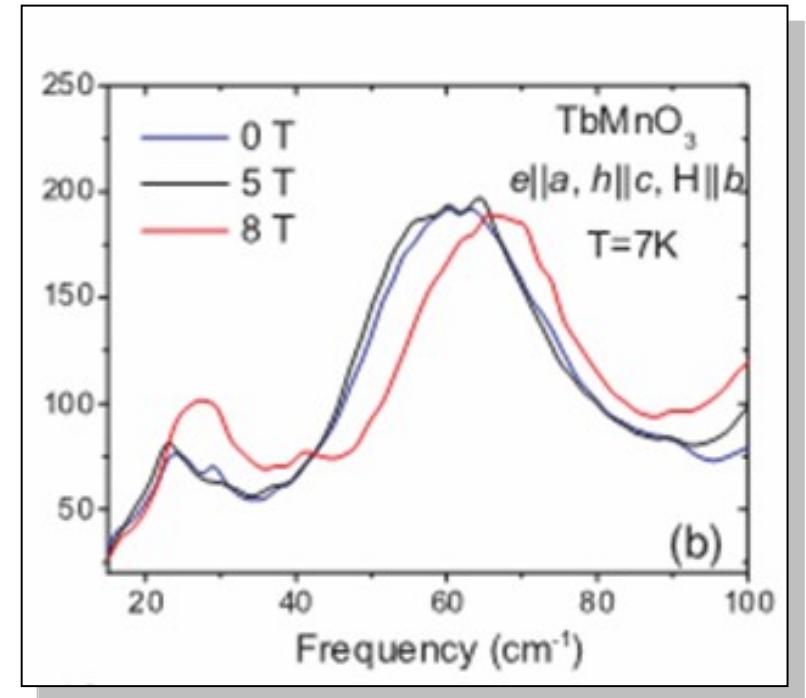
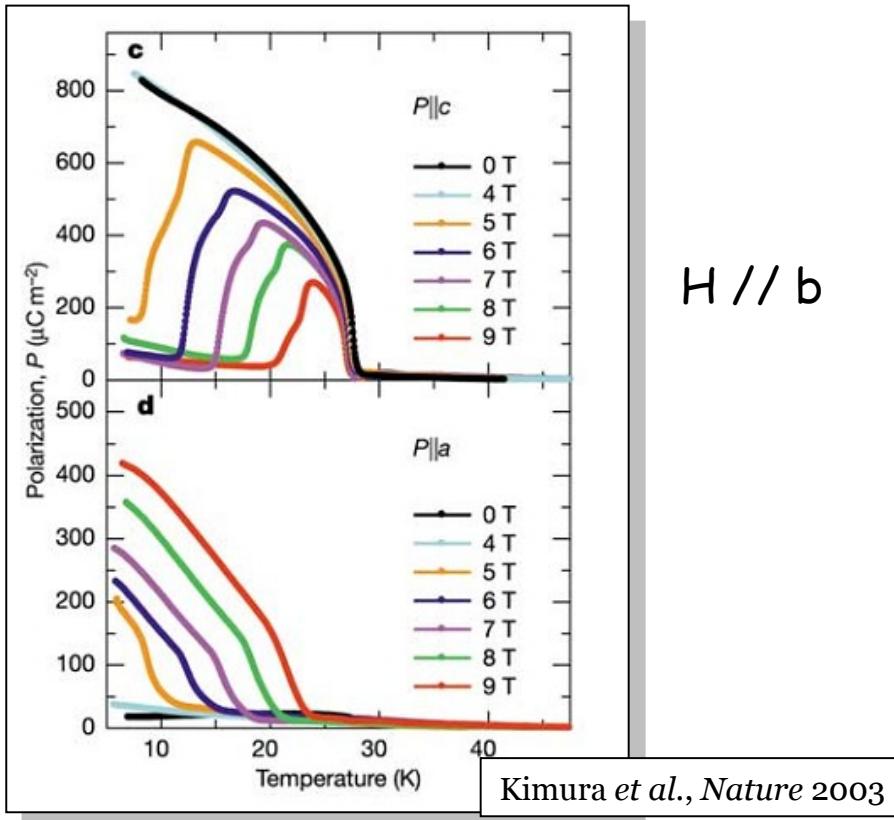
# The Electromagnon contributes to $\epsilon(0)$



Another confirmation that the electromagnon is excited by the electric field

# The electromagnon and polarization rotation

(Is the e-magnon related to a static coupling?)

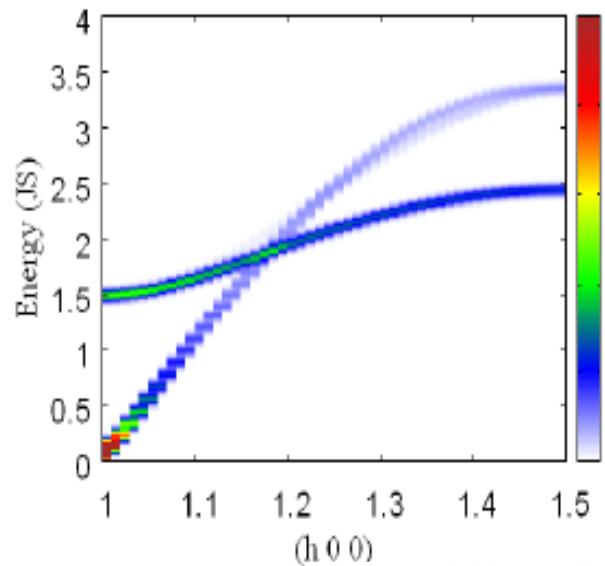


Valdés Aguilar et al., PRL, 2009.

The rotation of the polarization does not have a large effect on the electromagnon

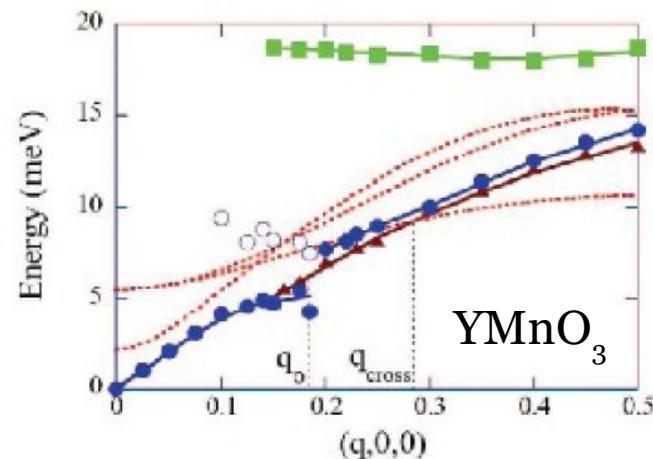
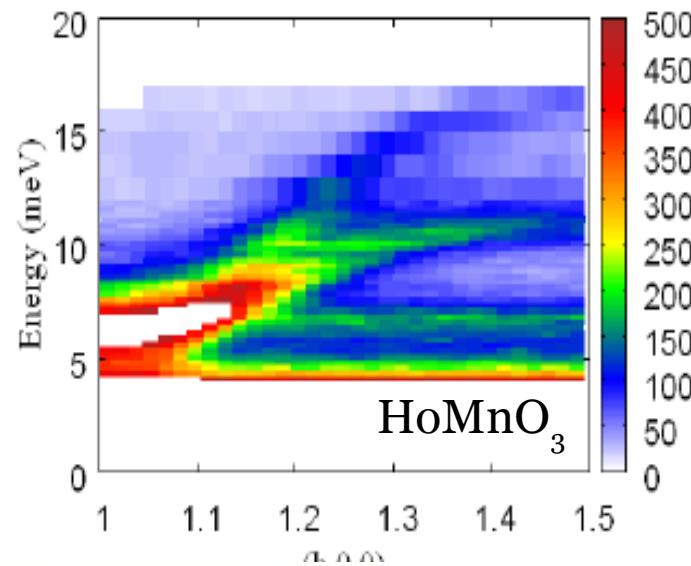
# Hybrid Mode – Neutrons

Calcul



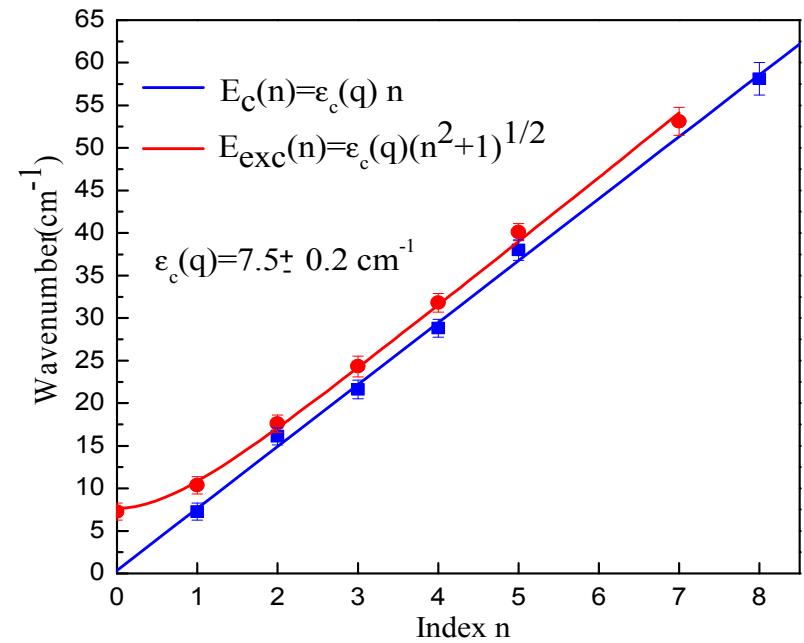
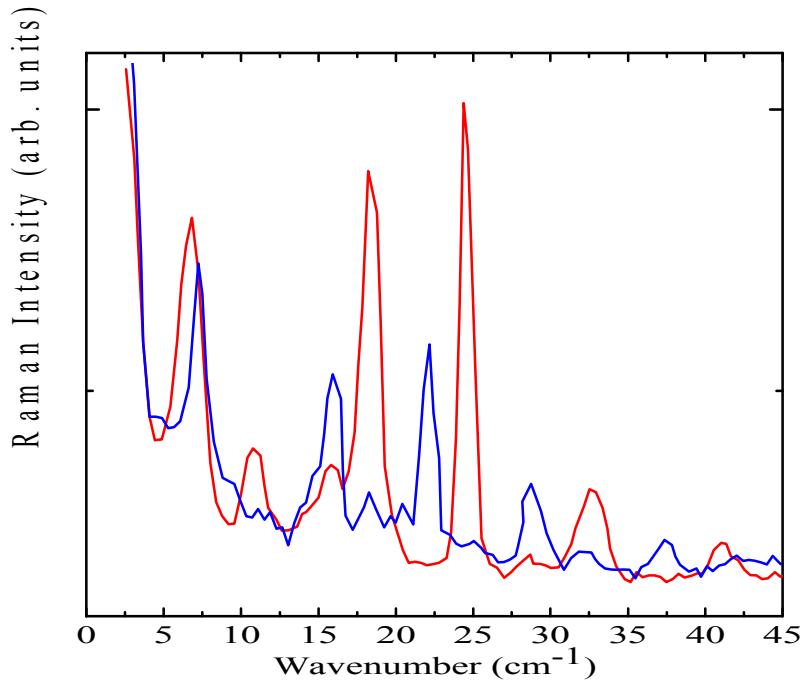
Mesures

HoMnO<sub>3</sub> T=1.5K



S. Petit, LLB, Saclay

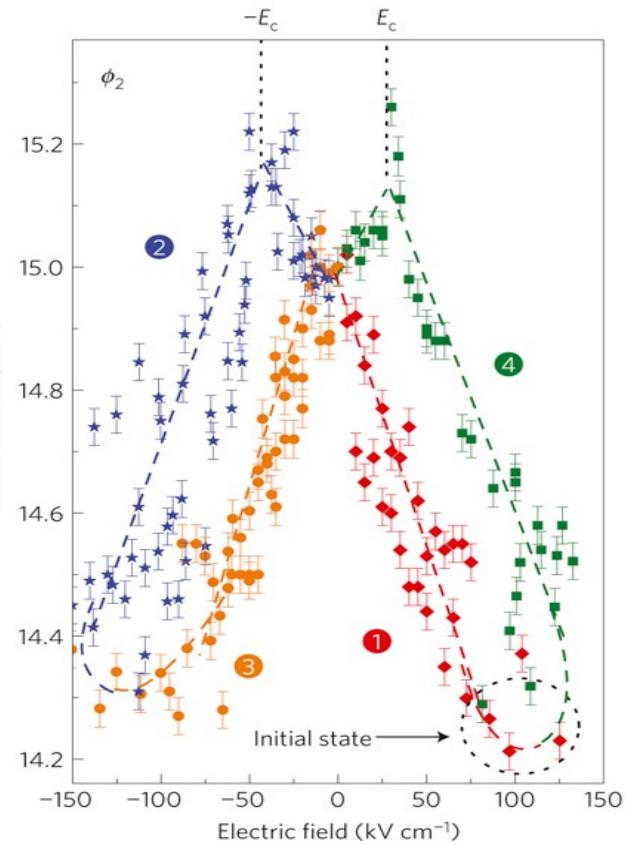
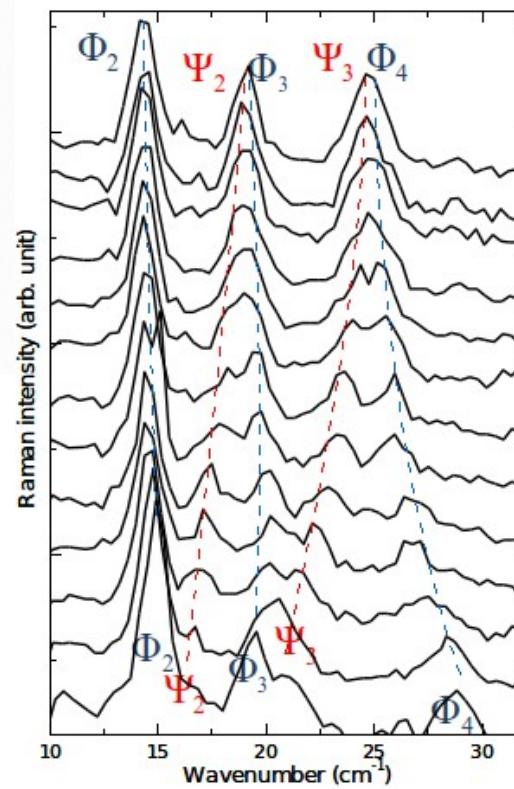
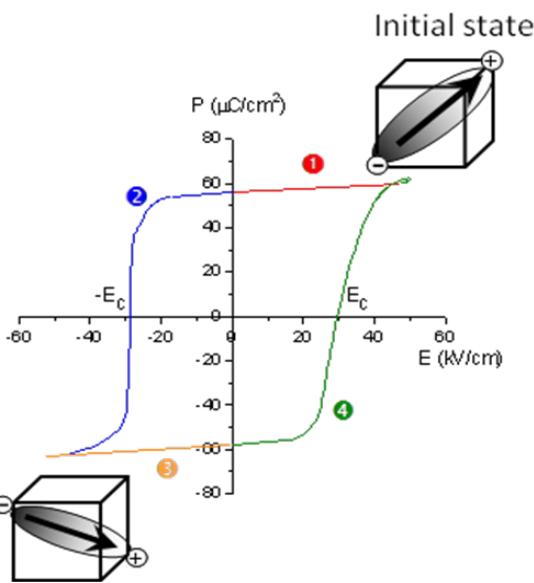
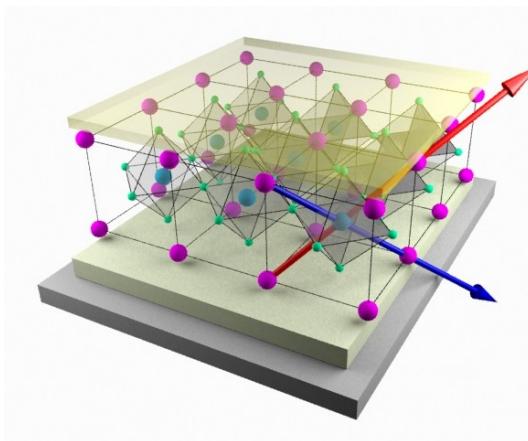
# (Electro?)mangons in $\text{BiFeO}_3$



- Two different spin waves
- Dispersion compatible with calculations for electromagnon

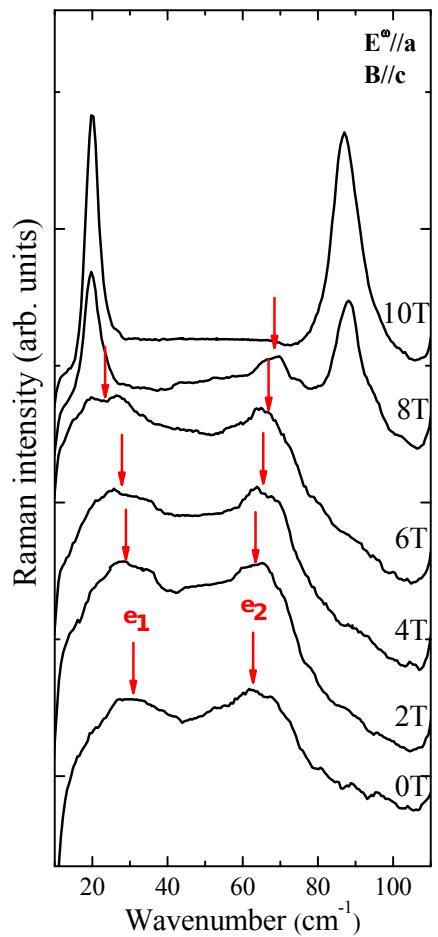
M. Cazayous et al., PRL 101, 03760 (2008)

# Electric Control of Spin Waves ( $\text{BiFeO}_3$ )



P. Rovillain et al. Nat. Mater. **9**, 979 (2010)

# Magnetic Control of the Electromagnon ( $\text{TbMnO}_3$ )



- In  $\text{TbMnO}_3$  under magnetic field
- Hybrid magnon-mode phonon
- De-hybridization of electromagnon

Rovillain et al. PRL **107**, 027202 (2011)

# Open questions

- ✓ Every good theorist can put together a Landau-like theory for the magneto-electric coupling, but what is its microscopic nature?
- ✓ What is the role of the lattice in the interaction?
- ✓ How important is ferroelasticity in the magneto-electric coupling?
- ✓ Do strong correlations have any saying in the matter?
- ✓ Is “magnetic” spatial symmetry breaking enough to create ferroelectricity?
- ✓ Where and what is the damn electromagnon?
- ✓ Is there a difference between static and dynamic coupling?